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# Size estimation of nano-particles from synchrotron X-ray data 

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MASTER'S THESIS
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## Preface

This is a Master's thesis in department of mathematics and natural sciences at the University of Stavanger, as part of the study program of Master of Science in Mathematics and Physics during the 2015 Fall and 2016 Spring semesters.

As the need of reliable methods for determining nanoparticle size increases, so the aim of this study was to analysis the size estimation by synchrotron X-ray diffraction. Mathematica was used as the main tool for analysing diffractogram.

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#### Abstract

Three samples of $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ powder were heated up from $380^{\circ} \mathrm{C}$ to $540^{\circ} \mathrm{C}$ and then cooled down to $380^{\circ} \mathrm{C}$ using an in-house made heat blower with a thermo-couple based temperature controller. During the heating process, synchrotron X-ray diffraction data including diffraction angles and intensities were obtained by diffractometer-setup found in the Swiss-Norwegian beamlines (SNBL- BM01A) at the ESRF. The diffractograms of each samples at different temperatures were studied by Mathematica. $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ powder procured at "Sigma-Aldrich" did not show any diffraction peak in the diffractograms, but the $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ powder obtained from crushing a large single crystal by pestle and mortar showed silver peaks in all five lowest reflections, 111, 200, 220, 311 and 222.

The sizes of the silver were estimated by Scherrer equation. The ways of defining peak width and the choice of Scherrer constants were found to be vital for the accuracy of size estimations. Peak width was defined by two methods, namely FWHM and integral breadth. For FWHM method, the sizes of silver grew typically from around $50 \pm 10 \mathrm{~nm}$ to $70 \pm 10 \mathrm{~nm}$ when temperature increased from $380^{\circ} \mathrm{C}$ to $432{ }^{\circ} \mathrm{C}$. For integral breadth method, the sizes of silver grew typically from around $60 \pm 10 \mathrm{~nm}$ to $80 \pm 10 \mathrm{~nm}$ when temperature increased from $380^{\circ} \mathrm{C}$ to 432 ${ }^{\circ} \mathrm{C}$.

For both methods, silver crystallites were grown almost linearly with increasing temperature in all reflections within the temperature range of $380^{\circ} \mathrm{C}$ to $432^{\circ} \mathrm{C}$. Then, silver crystallites stopped growing when temperature was over $432^{\circ} \mathrm{C}$ that is close to the phase transition temperature of silver sulphate, $427^{\circ} \mathrm{C}$. It is very likely that the phase transition of silver sulphate prohibits the growth of metallic silver.

The main instrumental broadening was coming from the finite pixel size of the detector and the energy dispersion of the synchrotron radiation. The uncertainties in FWHM were $0.04^{\circ} \pm$ $0.01^{\circ}$ to $0.10^{\circ} \pm 0.01^{\circ}$. The magnitude of the instrumental broadening due to the energy dispersion of the synchrotron radiation was $10^{-50}$ and was insignificant.


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## Abbreviations

ESRF European Synchrotron Radiation Facility

FCC Face Centered Cubic

FWHM Full Width at Half Maximum

MAE Mean Absolute Error

TDS Thermal Diffuse Scattering

TEM Transmission Electron Microscope

XRD X-Ray Diffraction

## Chapter 1

## Introduction

### 1.1 Background

Nanoparticles refer to particles with dimensions measured in nanometer (nm). In 2008, the International Organization for Standardization (ISO) defined a nanoparticle to be a discrete nanoobject that all three Cartesian dimensions have to be less than 100 nm [1]. Nanoparticles can be found in natural or can be made by human. During the last few decades, nanoparticles have become more important due to the practical applications in a variety of areas, including medicine, engineering, catalysis, and environmental remediation [1].

Therefore, the need of reliable methods for determining nanoparticle size has increased and this study was dedicated to improve the size estimation by X-ray diffraction.

## Nanoparticles to be studied

Silver nanoparticles have special optical, electronic and chemical properties, so they are used in many fields such as antibacterial application [2], catalysis [3, 4]. The study of silver nanoparticles has become more important, so silver nanoparticles were chosen to be studied.

In this research project, silver nanoparticles precipitating in silver sulphate matrix were studied. According to Larsen et al. [5], during the heating process of silver sulphate, metallic silver precipitated inside the silver sulphate matrix. The metallic silver will modify the physical properties associated with silver sulphate like ionic conduction property, so it is useful to understand the growth of silver precipitates in silver sulphate.

### 1.2 Method to study nanoparticles

Transmission electron microscopy (TEM) and X-ray diffraction (XRD) [6] are the most common methods to estimate the size of nanoparticles.

TEM is a direct observation method. The size of particles can be read directly. TEM is mostly used for studying the fine structure of crystallites. However, the samples have to be made into very thin layers so that they are observable in TEM. The preparations of the samples are difficult for some materials. In addition, some very small particles can be transparent for electrons [6]. Also, TEM provides information about the local structure, but the nanoparticles can have a distribution of sizes. So, the size estimations by TEM can have great error if there are a great variety of sizes in the nanoparticles [6, 7].

XRD method is the study of diffraction pattern from the crystallites. The size of the crystallites can be estimated from the information in the diffraction profiles, like peak breadth and peak intensity, by Scherrer equation. XRD can overcome the difficulties and limitations faced when using TEM. Firstly, XRD can be used with large powder, so preparations of very small samples are not necessary. Secondly, the recorded X-ray diffraction patterns come from a larger sample volume and thus XRD gives a volume averaged result. Therefore, XRD is more favorable than TEM in this study. However, it is worth to mention that there are some limitations in the use of Scherrer equation, which would be mentioned in Chapter 4.

### 1.3 Objectives

There are four main goals in this study:

- To examine the presence of silver in the silver sulphate in different samples.
- To study the growth of metallic silver in silver sulphate during the heating process.
- To compare the size estimations by FWHM and integral breadth method.
- To study the growth kinetics of silver


## Chapter 2

## Experiment

The experimental data were coming from research project done by Larsen et al. [5]. Below was given, for the sake of completeness, a resume of the experimental work. The majority of experimental data were analyzed by Mathematica, while some were by Excel.

## Experimental details

Three samples of silver sulphate were tested in this study. Sample 1 was some silver sulphate powder procured at "Sigma-Aldrich" and sample 3 was the same as sample 1 but having smaller grain sizes. Sample 2 was silver sulphate powder obtained from crushing a large single crystal by pestle and mortar. The single crystal was prepared according to a procedure given in research project done by Larsen et al. [5].

Each sample was filled in thin-walled glass capillaries with an outer diameter of 0.3 mm and mounted on a rotary axis with rotation speed $\omega$ at the diffractometer-setup found in the SwissNorwegian beamlines (SNBL- BM01A) at the ESRF. After that, the sample was heated up to $540^{\circ} \mathrm{C}$ from $380^{\circ} \mathrm{C}$, and then cooled down back to $380^{\circ} \mathrm{C}$ at a rate of $2^{\circ} \mathrm{C}$ per minute, excepting sample 2 was cooled down to $510^{\circ} \mathrm{C}$ only due to some technical errors in software communication setup. The temperature measurements were conducted using an in-house made heat blower with a thermo-couple based temperature controller. This was placed closed to the sample. The error of the thermo-couple was $\pm 0.5^{\circ} \mathrm{C}$.

During the heating process, collimated synchrotron beams of wavelength 0.6941 Å were produced by Swiss-Norwegian beamline A setup as shown in Figure 2.1 [8] with optics shown in Ta-


Figure 2.1: Beamline layout at the Swiss-Norwegian beamline A

Table 2.1: Optics for BM01A

| Optical elements | Mirror 1 | Double-crystal mono. | Mirror 2 |
| :---: | :---: | :---: | :---: |
| Distance from source | 25.9 m | 28.5 m | 30.8 m |
| Focusing type | Rh coated | Si (111) | Rh coated |
|  | Vert. collim. | Sagittally focusing | Vert. focusing |
| Beam size at sample | Nominally $0.5 \times 0.5 \mathrm{~mm}^{2}$ |  |  |
|  | FWHM focused |  |  |
| Spectral range | $6-22(30) \mathrm{keV}$ |  |  |
| Horizontal acceptance | 2 mrad |  |  |



Figure 2.2: Experimental diffraction geometry
ble 2.1 [8]. Diffraction pattern was produced as shown in Figure 2.2. The diffracted angles and diffracted intensities were recorded by the Dectris PILATUS2M Pixel detector with $0.172 \mathrm{~mm} \times$ 0.172 mm pixel size. As shown in Figure 2.2, the diffracted angles $2 \theta$ were the angles spread by each circle from the reference point of the powdered sample to the detector, while the recorded diffracted intensities were the azimuthal integration using the software package "Fit2D".


Figure 2.3: Diffractogram of powdered $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ at $380^{\circ} \mathrm{C}$

Then, diffractograms were plotted by Mathematica with the recorded diffraction angles and the diffracted intensities at different temperature respectively. Figure 2.3 shows an example of diffractogram of sample 1 at $380^{\circ} \mathrm{C}$. The diffraction peak profiles for all crystallites contained in the samples and all valid reflections of silver and silver sulphate can be found from the diffractogram.

In this study, the five lowest orders of silver peaks were tried to be found in each sample so as to examine the presence of silver throughout the heating process of silver sulphate. To do so, the diffraction angles were calculated with Bragg's law. Then, the diffractograms of that particular angles range were studied to identify the presence of silver peaks. If the diffractograms show silver peaks at the diffraction angle calculated, the presence of silver in the sample is proved. By studying the silver peak profiles, the size of the silver crystallites can be calculated with Scherrer equation.

## Chapter 3

## Formation of metallic silver

The work of Masciocchi and Parrish [9] stated that when powder of silver sulphate was heated up to $425{ }^{\circ} \mathrm{C}$, silver sulphate would transit from the low-temperature orthorhombic phase to hexagonal phase and would produce metallic silver. According to the research project done by Masciocchi and Parrish [9], the diffractograms showed that all silver sulphate was transited to hexagonal phase at $425^{\circ} \mathrm{C}$. Meanwhile, by looking at silver 111 reflection peak, it showed that the powder only contained a small volume of silver precipitate at $38.1^{\circ} \mathrm{C}$, and remained small until temperature was increased up to $350^{\circ} \mathrm{C}$. The silver precipitate grew rapidly at higher temperature, but Masciocchi and Parrish [9] did not mention about the change of volume of silver precipitate when temperature was heated up to temperature beyond $430^{\circ} \mathrm{C}$.

In this study, three samples of silver sulphate were tested. In each sample, the silver sulphate was heated up from $380^{\circ} \mathrm{C}$ to $540^{\circ} \mathrm{C}$ and then cooled down back to $380^{\circ} \mathrm{C}$ at a rate of $2^{\circ} \mathrm{C}$ per minute, excepting sample 2 which was cooled down to $510^{\circ} \mathrm{C}$ only because of the technical error occurred in the experiment. The objectives of this part of the experiment were:

- To test which samples would produce metallic silver when they were heated up.
- To look at the growth of metallic silver when the samples were heated up to temperature over the phase transition temperature of silver sulphate, which was $427^{\circ} \mathrm{C}$ in this study.


### 3.1 Diffraction angle

### 3.1.1 Bragg's law

The diffraction angle of the silver can be found by Bragg's law [10]:

$$
\begin{equation*}
2 d_{h k l} \sin \theta=\lambda \tag{3.1}
\end{equation*}
$$

where $h k l$ is the Miller index, $d_{h k l}$ is the inter-planar spacing at $h k l$ reflection, $\theta$ is the Bragg angle and $\lambda$ is the wavelength of X-ray used, which is $0.6941 \AA$ in this study. Since the metallic silver is FCC structure, the inter-planar spacing is [11]:

$$
\begin{equation*}
d_{h k l}=\frac{a_{\text {silver }}}{\sqrt{h^{2}+k^{2}+l^{2}}} \tag{3.2}
\end{equation*}
$$

where $a_{\text {silver }}$ is the lattice parameter of silver.

### 3.1.2 Extinction rule

Some of the reflections of silver can not be observed in diffractograms, due to the extinction rule of FCC structure. The corresponding extinction rule is: [12]:

- FCC structure extinction rule: all $h, k, l$ are mixed odd or even.

Therefore, the five lowest orders of non-extinct reflections are 111, 200, 220, 311 and 222. By putting $h k l$ equal to 111, 200, 220, 311 and 222 in Equation 3.2 respectively, $d_{h k l}$ of the five lowest orders of reflections were found and with these values, the diffraction angles were found by Equation 3.1. ${ }^{1}$

### 3.1.3 Lattice constant of silver

Since lattice constant changes with temperature [13, 14], to find an accurate diffraction angle of the silver, the lattice constants of silver at different temperatures have to be found.

[^0]In this study, data of silver's lattice constants at different temperatures from the Spreadborough and Christian's experiment [15] were collected. The data were used to interpolate the lattice constants of silver at different temperatures. The result is shown as Table 3.1

Table 3.1: Interpolation of silver lattice constant by linear regression and exponential regression, where a is the lattice constant in $\AA$ and T is the temperature in ${ }^{\circ} \mathrm{C}$

| Linear regression | Exponential Regression |
| :---: | :---: |
| $\mathrm{a}=0.00009^{*} \mathrm{~T}+4.083$ | $\mathrm{a}=4.0831^{*} \exp \left(0.00002^{*} \mathrm{~T}\right)$ |

The difference between the lattice constants generated by linear regression and exponential regression ${ }^{2}$ [16] were of magnitude smaller than 0.001 ( $2 \theta$, degree), which would not affect the identification of silver peaks. The silver peaks could be identified by both methods. However, overlapping of $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ peaks with Ag peaks could occur. The solution to this problem would be discussed in the coming section.

To compare the accuracy of linear and exponential regression, mean absolute error (MAE) [17] was used.

$$
\begin{equation*}
M A E=\frac{1}{n} \sum_{i=1}^{n}\left|f_{i}-y_{i}\right| \tag{3.3}
\end{equation*}
$$

where $n$ is the number of data, $f_{i}$ is the predicted value by regression and $y_{i}$ is the true value. As shown in Table 3.2, exponential regression on average, had a slightly lower mean absolute error than linear regression, so exponential regression is slightly better.

Table 3.2: Mean absolute errors of linear regression linear regression and exponential regression with the experimental diffraction angle

| Reflection | MAE of linear regression, <br> $2 \theta$ (degree) | MAE of exponential regression, <br> $2 \theta$ (degree) |
| :---: | :---: | :---: |
| 111 | 0.0102 | 0.0098 |
| 220 | 0.0113 | 0.0109 |
| 222 | 0.0163 | 0.0158 |
| 311 | 0.0140 | 0.0133 |
| 222 | 0.0209 | 0.0202 |
| Average | 0.0145 | 0.0140 |

[^1]One thing worth to mention is that for 220 reflection, the silver peaks should be identified with caution, as there was a side peak originating from $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ next to the silver peak that would affect the identification of silver peak as shown in Figure 3.1. So, addition code in Mathematica had to be added to avoid recording wrong diffraction angle and intensity. ${ }^{3}$


Figure 3.1: $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ peak nearby silver 220 reflection peak

## Interpolation of lattice constants

The data of lattice constants in the experiment were done by Spreadborough and Christian [15]. It contained values that were not within our experimental temperature range. So, the interpolation of lattice constants was performed again, with those values which were close to our experimental temperatures. These values were then compared with lattice constants that were interpolated by entire data set.

Table 3.3 shows that the mean absolute error on average became lower, if we interpolated with only lattice constants that were closed to our experimental temperature range, which was within $300^{\circ} \mathrm{C}$ to $600^{\circ} \mathrm{C}$. It implies that if no further experiment data can be provided, and interpolation of lattice constant is needed for future studies, then interpolation with those values, which are close to your experimental temperature range is recommended.

[^2]Table 3.3: Difference on interpolation of diffraction angle with data within $300^{\circ} \mathrm{C}$ to $600^{\circ} \mathrm{C}$ and interpolation with entire data set by exponential regression

| Reflection | MAE of exponential regression with data <br> within $300^{\circ} \mathrm{C}$ to $600^{\circ} \mathrm{C}, 2 \theta\left({ }^{\circ}\right)$ | MAE of exponential regression with <br> entire data set, $2 \theta\left({ }^{\circ}\right)$ |
| :---: | :---: | :---: |
| 111 | 0.00548 | 0.0098 |
| 220 | 0.00579 | 0.0109 |
| 222 | 0.01103 | 0.0158 |
| 311 | 0.00664 | 0.0133 |
| 222 | 0.00996 | 0.0202 |
| Average | 0.00778 | 0.0140 |

### 3.2 Examining the presence of silver in the samples

To test if the samples contained silver, the diffractograms of the three samples were examined by Mathematica. ${ }^{4}$

From the last section, the diffraction angles of silver 111 reflection were calculated. According to Appendix A, diffraction angles of silver 111 reflection were within the range of $16.65(2 \theta)$ to $16.85(2 \theta)$. So, the maximum peak intensities within the range of $16.65(2 \theta)$ to $16.85(2 \theta)$ were recorded to test if 111 reflection silver peaks were there.

### 3.2.1 Exponential model for decay of synchrotron beams

Since the synchrotron beams will decay exponentially and it implies that the intensity of the diffracted beams will decrease with time. This depends on the mode of operation of the synchrotron, which is the amount of electrons that are injected into the storage ring. The exponential decay equation is[18]:

$$
\begin{equation*}
\frac{d N}{d t}=-t_{1 / 2} N \tag{3.4}
\end{equation*}
$$

where $N$ is the quantity or the intensity in this study, $t_{1 / 2}$ is the half life which is 10 hours for the synchrotron beams being used and $t$ is the elapsed time. The solution to Equation 3.4 is [18]:

$$
\begin{equation*}
N=N_{0} e^{-t_{1 / 2} t} \tag{3.5}
\end{equation*}
$$

[^3]where $N_{0}$ is the initial quantity. By simple manipulation of Equation 3.5, the scaling factor $K$ is shown in Equation 3.6.
\[

$$
\begin{equation*}
K=\frac{N_{0}}{N}=\frac{1}{e^{-t_{1 / 2} t}} \tag{3.6}
\end{equation*}
$$

\]

The effect of the exponential decay can then be cancelled by multiply the recorded diffraction intensities by the scaling factor $K^{5}$.


Figure 3.2: Influence on peak intensities for silver 111 reflections without proper scaling

Figure 3.2 shows that the peaks intensities became higher after scaling and that means the actual diffracted intensities should be higher than the recorded intensities. Also, Figure 3.3 illustrates that when sample was heated up to $430{ }^{\circ} \mathrm{C}$, the peak intensities for the scaled data set remained at certain level, while the peak intensities for the non-scaled data set trended to decrease across the temperatures. This implies that the actual silver peak intensity will only increase when the heating process is below $430^{\circ} \mathrm{C}$, and remain more or less the same after $430^{\circ} \mathrm{C}$.

### 3.2.2 No silver formation in sample 1 and sample 3

With the scaling factor calculated from the previous section, the intensities of silver 111 reflections of the 3 different samples at different temperatures were calculated. Figure 3.4 shows that the maximum peak intensities within the range of $16.65(2 \theta)$ to $16.85(2 \theta)$ for sample 1 and sample 3 remained low across the temperature range of $380^{\circ} \mathrm{C}$ to $530^{\circ} \mathrm{C}$. Also, by taking a closer

[^4]

Figure 3.3: Decreasing in peak intensities for silver 111 reflection without proper scaling after reaching the phase transition temperature of silver sulphate


Figure 3.4: Intensities of silver 111 reflection of the 3 different samples along with temperature


Figure 3.5: Diffractogram to examining the presence of silver in the samples at temperature 530 ${ }^{\circ} \mathrm{C}$ in 111 reflection
look of the diffractogram at $530^{\circ} \mathrm{C}$ in Figure 3.5, it is observed there was no silver peak for sample 1 and sample3. These prove that there is no silver formation for sample 1 and sample 3 throughout the process.

### 3.2.3 Silver formation in sample 2

Figure 3.4 and Figure 3.5 indicate that silver peaks were found and that means silver was formed in sample 2. Figure 3.6 shows an example of diffractogram of sample 2 at $430^{\circ} \mathrm{C}$ that indicating the presence of silver peaks for the five lowest reflections. The reason of why only sample 2 had silver peaks might be explained by the difference in producing the silver sulphate powder. Since, sample 2 was made from one single silver sulphate crystal, metallic silver was precipitated inside the silver sulphate matrix, possibly by disproportionation. Metallic silver crystallites were embedded in sample 2 before the heating process and started to grow in size when they were heated up, while the sample 1 and sample 3 were coming from the commercial powder that did not contain any metallic silver. More details about the origin of the metallic silver would be discussed in Chapter 7.

From Figure 3.4, silver was formed at $380^{\circ} \mathrm{C}$, which was before the heating process was started. That means silver can be formed at temperature lower than the phase transition temperature of silver sulphate. This disagrees with Masciocchi and Parrish [9]'s saying that silver


Figure 3.6: Diffractogram of sample 2 at $530^{\circ} \mathrm{C}$ that showing the five lowest reflection peaks of silver
sulphate would transit from the low-temperature orthorhombic phase to hexagonal phase and produce metallic silver, because silver can be formed before the phase transition begins. Also, Figure 3.4 reveals that the peak intensity reached plateau after the phase transition. it is very likely that the phase transition of silver sulphate prohibits the growth of metallic silver rather than leads to the growth of silver. Further investigation is needed to understand the mechanism for the growth of metallic silver and some considerations would be given later in the thesis.

From Figure 3.4, it is shown that the peak intensity increased with the temperature and reached a plateau when temperature was increased to $430^{\circ} \mathrm{C}$ which was close to the phase transition temperature of the silver sulphate, $427^{\circ} \mathrm{C}$. However, the peak intensity does not represent the size of the silver. To study the growth of the silver size, Scherrer equation is needed and the size growth would be discussed in chapter 4.

## Chapter 4

## Size estimation of silver in sample 2

The size estimation of small crystallite powder from diffraction pattern is rooted from equation that describes the intensity of a diffracted peak from a small single parallelopipedon crystal [19]:

$$
\begin{equation*}
I=I_{e} F^{2} \frac{\sin ^{2}(\pi / \lambda)\left(\mathbf{s}-\mathbf{s}_{\mathbf{0}}\right) \cdot N_{1} \mathbf{a}_{\mathbf{1}}}{\sin ^{2}(\pi / \lambda)\left(\mathbf{s}-\mathbf{s}_{\mathbf{0}}\right) \cdot \mathbf{a}_{\mathbf{1}}} \frac{\sin ^{2}(\pi / \lambda)\left(\mathbf{s}-\mathbf{s}_{\mathbf{0}}\right) \cdot N_{2} \mathbf{a}_{\mathbf{2}}}{\sin ^{2}(\pi / \lambda)\left(\mathbf{s}-\mathbf{s}_{\mathbf{0}}\right) \cdot \mathbf{a}_{2}(\pi / \lambda)\left(\mathbf{s}-\mathbf{s}_{\mathbf{0}}\right) \cdot N_{3} \mathbf{a}_{\mathbf{3}}} \sin ^{2}(\pi / \lambda)\left(\mathbf{s}-\mathbf{s}_{\mathbf{0}}\right) \cdot \mathbf{a}_{\mathbf{3}} \tag{4.1}
\end{equation*}
$$

and [19],

$$
I_{e}=I_{0} \frac{e^{4}}{m^{2} c^{4} R^{2}} p
$$

where $I$ indicates the diffracted beam intensity, $I_{0}$ is the intensity of the polarized primary beam, $p$ is the polarization factor, for conventional X-ray source, $p$ is equal to $\frac{1+\cos ^{2} 2 \theta}{2}$ and for ideal synchrotron X -ray source, $p$ is equal to 1 in the plane [20]. $F$ represents the structure factor, $e$ and $m$ represent the charge and mass of electron, $c$ is the velocity of light, $\mathbf{s}_{\mathbf{0}}$ and $\mathbf{s}_{\mathbf{0}}$ are unit vectors that give the direction of the primary beam and the direction to the point of observation, $R$ is the distance from the crystal to the point of observation for the scattered beam and $N_{1,2,3}$, are the numbers of unit cells along the lattice directions $\mathbf{a}_{\mathbf{1 , 2 , 3}}$.

In general, if the crystallite size is large which is $N_{1,2,3}$ are large numbers, then each of the three quotients are non-zero only if the three Laue equations are closely satisfied and the diffraction peaks are sharp [19]. As shown in Figure 4.1, no size information can be told from the peak profile.

On the other hand, if the crystallite size is small where $N_{1,2,3}$ are small enough, the three
quotients are broadened as shown in Figure 4.2, and the breadth of the diffracted peak is inversely proportional to the size of the crystallite. So, the peak breadth can then be used to find the crystallite size with size range up to about $1000 \AA$ [19].


Figure 4.1: The function $\frac{\sin ^{2} N x}{\sin ^{2} x}$ for $N=500$


Figure 4.2: The function $\frac{\sin ^{2} N x}{\sin ^{2} x}$ for $N=5$

The size estimation depends on the definition of the peak breadth and the estimation of the crystallite shape factor[21] or the Scherrer constant. These two factors would be discussed in the following sections.

### 4.1 Scherrer equation

Scherrer came up with the idea of size estimation of crystallites from the diffraction profile, and the Scherrer equation is [22,23]:

$$
\begin{equation*}
D_{h k l}=\frac{K \lambda}{B_{h k l} \cos \theta} \tag{4.2}
\end{equation*}
$$

where $D_{h k l}$ is the crystallite size in the direction perpendicular to the lattice plane, $h k l$ represents the Miller index, $K$ is the Scherrer constant, $\lambda$ is the wavelength of the synchrotron radiation, $B_{h k l}$ represents the peak breadth of the diffracted peak and $\theta$ is the Bragg angle.

By studying the diffraction profile, the size of crystallite can then be told with known wavelength of radiation, and the corresponding diffraction angle.

### 4.1.1 Limitation to Scherrer equation

Warren [19] showed the detail derivation of Scherrer equation from Equations 4.1 and there are actually some assumptions during the derivation that limit the use of Scherrer equation. The assumptions include:

1. $N_{1}=N_{2}=N_{3}=N$
2. $\mathbf{a}_{1}, \mathbf{a}_{2}, \mathbf{a}_{3}$ are equal and orthogonal
3. Peak broadening is due to the small crystallite size only
4. Crystals are free from stains and faulting
5. $\frac{\sin ^{2} N x}{\sin ^{2} x}$ is approximated by Gaussian function $N^{2} e^{-(N x)^{2} / \pi}$

## Crystallites shape and size distribution

Assumption 1 implies that the powder crystallites are assumed to be cubic and all of the same size. Therefore, if the powder crystallites are not cubic in shape or have a size distribution, then Scherrer equation cannot be applied directly. Langford and Wilson [22] have summarized the modification needed for the choice of Scherrer constant if the crystallites were of other shape. The important of the choice of Scherrer constant would be discussed later in this Chapter. Wilson $[24,25]$ have considered the effect to the use of Scherrer equation when there was a size distribution and Langford and Wilson [22] have summarized the modification needed for the choice of Scherrer constant.

## Cubic symmetry

Assumption 2 means that the crystallites are assumed to be cubic symmetry. So, strictly speaking, Scherrer equation is only valid in cubic symmetry structure crystallites. The Scherrer equation holds for $A g$, since its structure is indeed cubic (FCC). Warren [19] stated that for other structure, it is not a bad approximation if for each $h k l$-reflection, the value of $D_{h k l}$ is assumed to be an average crystal dimension perpendicular to the reflecting planes.

## Broadening due to other factors

Assumptions 3 and 4 implies that diffraction peak profiles are only due to the small crystallite size. In reality, there are always others factors that will also contribute to the broadening, for instance, instrumental broadening, stains and faulting. These broadenings effect have to be eliminated before applying the Scherrer equation. Many researches [19, 26-29] have made effort on studying broadening due to other factors. These factors would be discussed in Chapter 5.

## Gaussian approximation

Assumption 5 implies that the diffraction peak profile is assumed to be a Gaussian function. The Gaussian fitting would be discussed in detail later in this Chapter. Langford and Wilson [22] have summarized the modification needed for diffraction peak profile having other functional form.

### 4.1.2 Scherrer constant

Before starting with the size estimation, the Scherrer constant and the peak breadth have to be defined. According to Langford and Wilson [22] Scherrer constant depends on three things:

- Definition of the breadth
- Crystallite shape
- Crystallite-size distribution

Different definitions of the breadth will give different values of Scherrer constant, the corresponding Scherrer constant value can be found from researches done by Langford and Wilson


Figure 4.3: Diffractogram showing FWHM of 200 reflection of silver peak at $398^{\circ} \mathrm{C}$
[22], Patterson [30, 31]. TEM can be used to find shape of crystallite and the corresponding Scherrer constant can be found from research done by Langford and Wilson [22].

In this study, there was no TEM indication on the crystallite shape and crystallite size distribution. So, crystallites were assumed to be cubic and the Scherrer constants for the 5 lowest reflections were found by Langford and Wilson's research [22]. Also, the effect of the crystallitesize distribution was omitted.

### 4.1.3 Define peak breadth by FWHM

The first way of defining the peak breadth was stated by Scherrer [23], that is the full width of the diffracted maximum at the height, which is half of the background and the maximum peak intensity (FWHM). Figure 4.3 shows an example of FWHM of 200 reflection of silver peak at 398 ${ }^{\circ} \mathrm{C}$.

## Background intensity

There was some background intensities, for example TDS, that would contribute to the silver peak intensity. So, as a first step to find the FWHM, the background intensities were cancelled by calculating the average value of recorded background intensities as shown in Figure 4.4.


Figure 4.4: Diffractograms showing the silver 200 reflection peak at $39{ }^{\circ} \mathrm{C}$ with and without background intensities

## Gaussian fitting

Due to the finite pixel size of the detector, the FWHM could not be found from the diffractogram directly. As an example, the 200 reflection peak at $398^{\circ} \mathrm{C}$ is shown in Figure 4.5 . The peak profile was composed of only few data points, and the FWHM could not be recorded directly.

The Gaussian fitting based on the data points that were recorded, and the Gaussian Equation is [32]:

$$
\begin{equation*}
f(x)=A e^{-\frac{(x-\mu)^{2}}{\sigma^{2}}} \tag{4.3}
\end{equation*}
$$

where $A$ is the height of the peak, $\mu$ is the position of the center of the peak and $\sigma$ is the standard derivation.

It is worthwhile to mention that in Mathematica, the parameters $A, \mu, \sigma$ have to be set as some reasonable pre-set values so that the Gaussian fitting can function properly ${ }^{1}$. The guesses can be made by observing the range of the diffraction angle and diffraction intensities for that particular reflection.

Another thing that is needed to pay attention to is the Gaussian fittings for 111 reflection and 220 reflections because there were side peaks nearby the silver peaks. So, those values have to be excluded from the fitting by chosing fitting range carefully ${ }^{1}$.

[^5]

Figure 4.5: Recorded diffraction angles and intensities within the range of 200 reflection at 398 ${ }^{\circ} \mathrm{C}$

One of the examples of Gaussian fitting is shown in Figure 4.6. The peak profile is now completed and the FWHM can then be recorded from the Gaussian fitting curve by the following equation [33]:

$$
\begin{equation*}
\mathrm{FWHM}=2 \sqrt{2 \ln 2} \sigma \tag{4.4}
\end{equation*}
$$

Figure 4.7 shows the calculated FWHM for all five lowest reflections of silver.

## Size estimation

With the calculated FWHM and the corresponding Scherrer constants from Langford and Wilson's research [22], the sizes of silver were estimated by Scherrer Equation 4.2. The results of the size estimations of silver for the 111, 200, 220, 311 and 222 reflections are shown in Figure 4.8.

Since the sample was heated up to $540^{\circ} \mathrm{C}$ and then cooled back to $510^{\circ} \mathrm{C}$, the plots were somehow overlapping in the tail. To make the plots clearer, the x -axis was changed to elapsed time in minutes and y-axis to nano-meter, as shown in Figure 4.9.


Figure 4.6: Gaussian fitting of silver 200 reflection at $398^{\circ} \mathrm{C}$


Figure 4.7: Calculated FWHM for all five reflections


Figure 4.8: The size estimations of silver for the five lowest reflections by defining peak breadth by FWHM


Figure 4.9: The size estimations of silver for the five lowest reflections by defining peak breadth by FWHM with x -axis as elapsed time in minutes and y -axis in nano-meter

### 4.1.4 Define peak breadth by integral breadth

The second way of defining the peak breadth was first introduced by Laue [34] (see also Stokes and Wilson [35] and Jones [36]), the peak breadth or the integral breadth is defined as the total area under the diffraction peak profile and divided by the peak intensity.

## Background and Gaussian fitting

The background intensity was cancelled with the same method as previous section shown in Figure 4.4. Also, the same Gaussian fitting was performed as the previous section shown in Figure 4.6.

## Size estimation

The total area under the silver peak profile generated by Gaussian fitting, as shown in Figure 4.10, was calculated by Mathematica. Then, the areas were divided by the peak intensity so as to find the integral breath. Figure 4.11 shows the calculated integral breadth for all five reflections. Together with the corresponding Scherrer constants from Langford and Wilson's research [22], the result of size estimations was shown in Figure 4.12.


Figure 4.10: Diffractogram showing peak intensity and the total area under diffraction peak profile of 200 reflection of silver peak at $398^{\circ} \mathrm{C}$


Figure 4.11: Calculated integral breadth for all five lowest reflections of silver


Figure 4.12: The size estimations of silver for the five lowest reflections by defining peak breadth by integral breadth and x -axis as elapsed time in minutes

### 4.2 Silver growth

From Figure 4.9 and 4.12 in the last section, the silver growth shows three features:

- Silver crystallites were growing almost linearly with temperature in all reflections.
- The size of silver crystallites was smallest in the beginning at $380^{\circ} \mathrm{C}$ and became the largest at around $432^{\circ} \mathrm{C}$ in all reflections.
- The silver crystallites stopped growing after reaching $432^{\circ} \mathrm{C}$ in all reflections
and Table 4.1 summarizes the size estimations of silver by FWHM method and integral breadth method.

Table 4.1: Summary on silver growth estimated by FWHM and integral breadth method

|  |  | Reflections |  |  |  |  |
| ---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Temperature | 111 | 200 | 220 | 311 | 222 |
| Size estimated by | $380^{\circ} \mathrm{C}$ | 58.40 | 59.64 | 50.00 | 38.84 | 60.87 |
| FWHM (nm) | $432^{\circ} \mathrm{C}$ | 72.82 | 72.34 | 72.45 | 75.06 | 79.00 |
| Size estimated by | $380^{\circ} \mathrm{C}$ | 74.08 | 63.25 | 59.74 | 45.64 | 77.36 |
| Integral Breadth (nm) | $432^{\circ} \mathrm{C}$ | 92.37 | 76.71 | 86.57 | 88.19 | 100.22 |

## Validity of Scherrer equation

According to Holzwarth and Gibson [37], the Scherrer equation can only be applied for size estimation up to about 100 nm to 200 nm , while the sizes of silver estimated by both methods were within the range of 50 nm to 100 nm . This implies that the size estimations with Scherrer equation are valid.

## Growth of metallic silver

Figures 4.9 and 4.12 show that silver sizes were growing during the heating process of silver sulphate. One of the suggestions to the growth of silver is the disproportionation theorem and more details would be discussed in Chapter 7.

## Stopping of metallic silver growth

With reference to Figures 4.9 and 4.12 , the silver stopped growing at $432^{\circ} \mathrm{C}$ which was close to the phase transition temperature which is $427^{\circ} \mathrm{C}$. This supports the observation from Chapter 3, which is the phase transition of silver sulphate prohibits the growth of metallic silver rather than leads to the growth of silver, because the silver crystallites grew before the phase transition and stopped at temperature near to the phase transition. To test the hypothesis, see Chapter 7 for further investigation suggested.

Another suggestion to the growth stop is coming from the saturation of metallic silver. According to Klimovich et al. [38], when the standard free energy change of the chemical reaction is non-negative, the reaction will stop. The metallic silver grow because it can reduce the standard free energy of the silver sulphate matrix. However, upon some levels, further growing of metallic silver can no longer reduce the free energy of the system. Therefore, the growth stops. Another effect could be due to some sort of site saturation. Further investigation is needed to test this hypothesis.

### 4.3 Uncertainty in size estimation

The uncertainty in size estimation can be expressed by Equation 4.5 [39].

$$
\begin{equation*}
\delta D_{h k l}=\sqrt{\left(\frac{\partial D}{\partial K} \delta K\right)^{2}+\left(\frac{\partial D}{\partial \lambda} \delta \lambda\right)^{2}+\left(\frac{\partial D}{\partial B_{h k l}} \delta B_{h k l}\right)^{2}+\left(\frac{\partial D}{\partial \theta} \delta \theta\right)^{2}} \tag{4.5}
\end{equation*}
$$

By performing the partial derivatives, Equation 4.5 becomes:

$$
\begin{equation*}
\delta D_{h k l}=\sqrt{\left(\frac{\lambda}{B_{h k l} \cos \theta} \delta K\right)^{2}+\left(\frac{K}{B_{h k l} \cos \theta} \delta \lambda\right)^{2}+\left(-\frac{K \lambda}{B_{h k l}^{2} \cos \theta} \delta B_{h k l}\right)^{2}+\left(\frac{K \lambda \sin \theta}{B_{h k l} \cos ^{2} \theta} \delta \theta\right)^{2}} \tag{4.6}
\end{equation*}
$$

## Uncertainty in Scherrer constant $\delta K$

The shapes of the silver crystallites were unknown, and the size estimations of silver crystallites were based on the Scherrer constant values for cubic shape crystallites. So, the uncertainty would be the maximum difference between the Scherrer constant of cubic shape with all other shapes listed in Table 4.2 and Table 4.3 [22]. $\delta K_{w}$ represents the uncertainty in Scherrer constant
for FWHM method and $\delta K_{\beta}$ represents the uncertainty in Scherrer constant for integral breadth method.

Table 4.2: Scherrer constants used in FWHM method for different shapes of silver crystallites and uncertainty of Scherrer constants, $\delta K_{w}$ that is the maximum difference between the Scherrer constant of cubic shape with all other shapes

|  | Scherrer constants for different shape of crystallites, $K_{w}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Reflection | Cubic | Spherical | Tetrahedral | Octahedral | $\delta K_{w}$ |
| 111 | 0.8551 | 0.9400 | 0.8894 | 0.9354 | 0.0849 |
| 200 | 0.8859 | 0.9400 | 1.0220 | 0.8151 | 0.2069 |
| 220 | 0.8340 | 0.9400 | 0.7262 | 0.8613 | 0.2138 |
| 311 | 0.9082 | 0.9400 | 0.9289 | 0.8863 | 0.0426 |
| 222 | 0.8551 | 0.9400 | 0.8894 | 0.9354 | 0.0849 |

Table 4.3: Scherrer constants used in integral breadth method for different shapes of silver crystallites and uncertainty of Scherrer constants, $\delta K_{\beta}$ that is the maximum difference between the Scherrer constant of cubic shape with all other shapes

| Scherrer constants for different shape of crystallites, $K_{\beta}$ |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Reflection | Cubic | Spherical | Tetrahedral | Octahedral | $\delta K_{\beta}$ |
| 111 | 1.1547 | 1.0747 | 0.8894 | 1.1438 | 0.2544 |
| 200 | 1.0000 | 1.0747 | 1.3867 | 1.1006 | 0.3120 |
| 220 | 1.0607 | 1.0747 | 0.9806 | 1.0377 | 0.0941 |
| 311 | 1.1359 | 1.0747 | 1.3156 | 1.1211 | 0.2409 |
| 222 | 1.1547 | 1.0747 | 1.2009 | 1.1438 | 0.1262 |

## Uncertainty in wavelength $\delta \lambda$

The uncertainty in wavelength was mainly coming from energy dispersion of the incoming synchrotron radiation. $\delta \lambda$ is $0.694 \times 10^{-15}$ and detail calculation is shown in Chapter 5 .

## Uncertainty in peak breadth $\delta B_{h k l}$

The uncertainty in peak breadth was coming from the Gaussian fitting. The $95 \%$ confidence interval of Gaussian fitting was found by Mathemcatica function and was used to find the uncertainty in peak breadth. Figure 4.13 shows an example of $95 \%$ confidence interval of Gaussian fitting of 200 reflection at $398^{\circ} \mathrm{C}$. The uncertainty in peak breadth was calculated by the
difference between the peak breadth of Gaussian fitting model and the peak breadth of $95 \%$ confidence interval. The value of $\delta B_{h k l}$ is taken as:

$$
\delta B_{h k l}=\frac{\left|B_{G}-B_{u}\right|+\left|B_{G}-B_{l}\right|}{2}
$$

where $B_{G}$ is the peak breadth of Gaussian fitting model, $B_{u}$ is peak breadth of the upper limit of $95 \%$ confidence interval and $B_{l}$ is peak breadth of the lower limit of $95 \%$ confidence interval.


Figure 4.13: $95 \%$ confidence interval of Gaussian fitting of silver 200 reflection at $398{ }^{\circ} \mathrm{C}$

Table 4.4 and 4.5 show an example of $\delta B_{h k l}$ of FWHM method and integral breadth method at $380^{\circ} \mathrm{C}$ and $432^{\circ} \mathrm{C}$ respectively.

Table 4.4: $\delta B_{h k l}$ of FWHM method and integral breadth method at $380^{\circ} \mathrm{C}$

| Reflection | $B_{F W H M}$ | $\delta B_{F W H M}$ | $B_{\text {Integralbreadth }}$ | $\delta B_{\text {Integralbreadth }}$ |
| :---: | :---: | :---: | :---: | :---: |
| 111 | 0.0589 | 0.0022 | 0.0627 | 0.0018 |
| 200 | 0.0599 | 0.0023 | 0.0638 | 0.0021 |
| 220 | 0.0683 | 0.0047 | 0.0727 | 0.0031 |
| 311 | 0.0968 | 0.0022 | 0.1031 | 0.0020 |
| 222 | 0.0584 | 0.0035 | 0.0621 | 0.0033 |

## Uncertainty in diffraction angle $\delta \theta$

The dominating part of uncertainty in diffraction angle was coming from finite pixel size of the detector and $\delta \theta$ is 0.35 . The detail calculated can be found in Chapter 5.

Table 4.5: $\delta B_{h k l}$ of FWHM method and integral breadth method at $432^{\circ} \mathrm{C}$

| Reflection | $B_{F W H M}$ | $\delta B_{F W H M}$ | $B_{\text {Integral breadth }}$ | $\delta B_{\text {Integral breadth }}$ |
| :---: | :---: | :---: | :---: | :---: |
| 111 | 0.0475 | 0.0012 | 0.0506 | 0.0011 |
| 200 | 0.0495 | 0.0012 | 0.0527 | 0.0011 |
| 220 | 0.0480 | 0.0018 | 0.0510 | 0.0013 |
| 311 | 0.0557 | 0.0022 | 0.0593 | 0.0020 |
| 222 | 0.0458 | 0.0009 | 0.0488 | 0.0009 |

## The uncertainty of size estimation

The uncertainties of size estimations $\delta D_{h k l}$ were calculated by putting all the calculated values above into Equation 4.6. Table 4.6 summarizes the results at $380^{\circ} \mathrm{C}$ and $432{ }^{\circ} \mathrm{C}$. Figure 4.14 and 4.15 show the uncertainties in size estimations for the five lowest reflections of silver throughout the entire experimental range.

Table 4.6: Summary on uncertainties in estimated silver sizes by FWHM and integral breadth method

|  |  | Reflections |  |  |  |  |
| ---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Temperature | 111 | 200 | 220 | 311 | 222 |
| Size estimated by | $380^{\circ} \mathrm{C}$ | $58 \pm 6$ | $60 \pm 14$ | $50 \pm 13$ | $39 \pm 2$ | $61 \pm 7$ |
| FWHM (nm) | $432^{\circ} \mathrm{C}$ | $73 \pm 7$ | $72 \pm 17$ | $72 \pm 18$ | $75 \pm 4$ | $79 \pm 8$ |
| Size estimated by | $380^{\circ} \mathrm{C}$ | $74 \pm 16$ | $63 \pm 20$ | $60 \pm 6$ | $46 \pm 10$ | $77 \pm 9$ |
| Integral Breadth (nm) | $432^{\circ} \mathrm{C}$ | $92 \pm 20$ | $77 \pm 24$ | $87 \pm 8$ | $88 \pm 17$ | $100 \pm 10$ |

### 4.4 Comparison on FWHM and integral breadth

## Silver growth pattern

According to Figure 4.16, the growth patterns of silver crystallites were the same in all reflections. For each reflection, the estimated sizes of silver by both methods grew linearly in between temperature range of $380^{\circ} \mathrm{C}$ to $432^{\circ} \mathrm{C}$ and stopped at temperature around $432^{\circ} \mathrm{C}$. This implies that both methods are for sure studying the same crystallites as their estimated sizes have the same pattern.


Figure 4.14: The uncertainties in size estimations of silver for the five lowest reflections by defining peak breadth by FWHM


Figure 4.15: The uncertainties in size estimations of silver for the five lowest reflections by defining peak breadth by integral breadth

## Estimated size difference

Figure 4.16 shows that the estimated sizes between FWHM method and integral breadth method have some difference throughout the temperature range of the experiment in all five reflections. This is possibly because of the difference in defining peak breath and the values of Scherrer constant that were used. However, this does not mean that one of the methods is more accurate than the other. To compare the accuracy of size estimations from these two methods in this study, further investigation is needed to examine the true size of silver.

## Silver growth rate

The average difference between size estimations by FWHM and integral breadth were calculated. This average value was added to the sizes estimated by FWHM, so that the two plots reached almost the same level, for better comparison between the estimated growth rates.

As shown in Figure 4.17, the integral breadth method showed a faster growth rate in between temperature range of $380^{\circ} \mathrm{C}$ to $432^{\circ} \mathrm{C}$ for all five reflections. This is possibly due to the difference in defining the integral breadth. While, the growth rate reach almost zero for temperature higher than $432^{\circ} \mathrm{C}$ for all five reflections.


Figure 4.16: Difference on size estimations of silver by defining peak breadth by FWHM and integral breadth


Figure 4.17: Comparing the estimated sizes growth rates of silver by FWHM and integral breadth method by shifting up the plot of sizes estimated by FWHM method

### 4.5 The important of Scherrer constant on size estimation

From Figure 4.18, it is known that there was not much difference in size estimations for silver 311 reflection by assuming spherical crystallite where Scherrer constant equal to 0.94 ; and assuming cubic crystallite where numerical value of Scherrer constant was found from the Langford and Wilson's [22] research.


Figure 4.18: Difference on size estimations for 311 reflection by assuming spherical crystallite where Scherrer constant equal to 0.94 ; and assuming cubic crystallite where Scherrer constant equal to 0.9082

However, Figure 4.19, 4.20, 4.21, 4.22 illustrate an obvious difference in size estimations for 111, 200, 200, 222 reflections by assuming spherical crystallite and assuming cubic crystallite. This implies that the size estimations can be quite different from real size if crystallite shape is unknown. Therefore, TEM-indication is necessary for an accurate size estimation.


Figure 4.19: Difference on size estimations for 111 reflection by assuming spherical crystallite where Scherrer constant equal to 0.94 ; and assuming cubic crystallite where Scherrer constant equal to 0.8551


Figure 4.20: Difference on size estimations for 200 reflection by assuming spherical crystallite where Scherrer constant equal to 0.94 ; and assuming cubic crystallite where Scherrer constant equal to 0.8859


Figure 4.21: Difference on size estimations for 220 reflection by assuming spherical crystallite where Scherrer constant equal to 0.94 ; and assuming cubic crystallite where Scherrer constant equal to 0.8340


Figure 4.22: Difference on size estimations for 222 reflection by assuming spherical crystallite where Scherrer constant equal to 0.94 ; and assuming cubic crystallite where Scherrer constant equal to 0.8551

## Chapter 5

## Factors leading to the broadening of diffraction peak

As mentioned in the previous chapters, the particle size affects the diffraction peak profile, but this is not the only factor. Many researches [27, 29, 37, 40-42] have stated that there are other factors that will affect the diffraction peak profile. To get a more accurate size estimation of the crystallite, these kinds of broadening have to be studied and eliminated.

### 5.1 Instrumental broadening

There are many types of instrumental broadening, such as slit widths, imperfect focusing, beam size, penetration into the sample[26], etc. Depending on the experimental setup, different instrumental broadenings should be considered.

In this study, Dectris PILATUS2M Pixel detector and double crystal monochromator setup were used. The main instrumental broadening was coming from the finite pixel size of the detector and the energy dispersion of the synchrotron radiation.

### 5.1.1 Finite pixel size of the detector

According to Bragg's law (Equation 3.1), the diffraction angle is inversely proportional to interplanar distance $d$ and $d$ is proportional to temperature. So, diffraction angles should decrease
with increasing temperatures. However, the experimental diffraction angles were found to be


Figure 5.1: Theoretical and experiment values of diffraction angle of 200 reflection
decreased in a zig-zag way and Figure 5.1 shows an example from silver 200 reflection that the diffraction angles did not decrease linearly as predicted theoretically.


Figure 5.2: Diffractogram of silver 200 reflection at $388^{\circ} \mathrm{C}, 390^{\circ} \mathrm{C}, 392^{\circ} \mathrm{C}$ and $394^{\circ} \mathrm{C}$

The reason to these results is due to the finite pixel size of the detector. Figure 5.2 shows that for 200 reflection, when temperature increased, the diffraction angle decreased. However, due to the finite pixel size of the detector, the peak position remain the same from $388^{\circ} \mathrm{C}$ to $392^{\circ} \mathrm{C}$ and then jumped to other position at $394^{\circ} \mathrm{C}$ and made the zig-zag curve shown in Figure 5.1.

## Magnitude of error on size estimations of silver

The uncertainties in recorded diffraction angles due to finite pixel size of the detector were found by Equation 5.1.

$$
\begin{equation*}
\delta \theta=\tan ^{-1}\left(\frac{l}{d}\right) \tag{5.1}
\end{equation*}
$$

where,

$$
l=\sqrt{\left(\frac{w}{2}\right)^{2}+\left(\frac{w}{2}\right)^{2}}
$$

where $\delta \theta$ is the uncertainty in recorded diffraction angle due to finite pixel size of the detector, $l$ is the maximum possible error of the diffracted position measured, $d$ is the sample to detector distance and $w$ is the width of the pixel.

The pixel size of the detector was $0.172 \mathrm{~mm} \times 0.172 \mathrm{~mm}$, so $w$ was 0.172 mm and sample to detector distance of the experimental setup $d$ was 20 cm . Therefore, uncertainty in recorded diffraction angle or the uncertainty in measured FWHM was of the magnitude $10^{-2}$.

With reference to Chapter 4, the broadenings due to the samples were within the range of 4 to $10 \times 10^{-2}$ (degree), and the magnitude of the instrumental broadening due to finite pixel size of the detector was $10^{-2}$ (degree). This implies that instrumental broadening due to finite pixel size of the detector is probably the dominating instrumental broadening factor and the uncertainties in FWHM were $0.04 \pm 0.01$ (degree) to $0.10 \pm 0.01$ (degree).

### 5.1.2 Energy dispersion of the incoming synchrotron radiation

The double crystal monochromator setup was used to select the required energy of synchrotron radiation. The maximum possible error during the energy selection is:

$$
\begin{equation*}
\frac{\Delta E}{E}=\frac{\Delta \lambda}{\lambda}=10^{-4} \tag{5.2}
\end{equation*}
$$

where $E$ is the energy of the synchrotron radiation and $\lambda$ is the wavelength of the synchrotron radiation. By differentiate Equation 3.1 with $d_{h k l}$ fixed,

$$
\begin{equation*}
2 d_{h k l} \cos \theta=\frac{d \lambda}{d \theta} \tag{5.3}
\end{equation*}
$$

and by simple manipulation, Equation 5.3 become,

$$
\begin{equation*}
d \theta=\frac{d \lambda}{2 d_{h k l} \cos \theta} \tag{5.4}
\end{equation*}
$$

by substituting Equation 3.2 and 5.2 into Equation 5.4,

$$
\begin{equation*}
d \theta=\frac{\lambda \times 10^{-4}}{2 \frac{a}{\sqrt{h^{2}+k^{2}+l^{2}}} \cos \theta} \tag{5.5}
\end{equation*}
$$

With reference to diffraction angles of silver at Appendix A, the maximum possible errors for the 5 lowest reflections due to energy dispersion of sample 2 at $380^{\circ} \mathrm{C}$ were shown in Table 5.1 ${ }^{1}$.

Table 5.1: Maximum Instrumental broadening due to energy dispersion of synchrotron radiation

|  | Reflections, 20 (degree) |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Temperature | 111 | 200 | 220 | 311 | 222 |
| $380^{\circ} \mathrm{C}$ | $2.89 \times 10^{-5}$ | $3.33 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |

Table 5.1 and Appendix C show that the maximum possible errors due to energy dispersion were of magnitude $10^{-5}$ (degree), while the magnitude of the instrumental broadening due to the sample was $10^{-2}$ (degree). This implies that instrumental broadening due to finite pixel size of the detector is negligible.

To double check if the errors affect the results, Equation 5.6 was used to calculate the difference in size estimations with and without eliminating the energy dispersion errors. Figure 5.3 illustrates that, the size estimations of silver in sample 2111 reflection were almost the same for estimations with or without eliminating the energy dispersion broadening. This proves that the broadening is negligible.

### 5.1.3 Method to remove instrumental broadening

It is worth to mention the general method to remove the instrumental broadening. Warren [19] stated that, by running a standard peak using a large powder size sample, all particle size broadening is eliminated. Then, the recorded diffraction peak profile will be mainly contributed by

[^6]

Figure 5.3: Size estimations of silver in sample 2111 reflection with and without eliminating the energy dispersion broadening
instrumental broadening. So, the instrumental broadening can be removed by using Equation 5.6 or $5.7[19,28]$ depending on the shape of the peak profile.

$$
\begin{equation*}
\text { Gaussian shape : } B^{2}(h)=B^{2}(g)+B^{2}(f) \tag{5.6}
\end{equation*}
$$

$$
\begin{equation*}
\text { Cauchy shape : } B(h)=B(g)+B(f) \tag{5.7}
\end{equation*}
$$

where $B(g)$ is the half maximum breadth due to instrumental broadening, $B(f)$ is the half maximum breadth due to particle size and $B(h)$ is the half maximum breadth due to both instrumental broadening and particle size.

### 5.2 Internal elastic strains

According to Hammond [26], the elastic strains can be separated into macro-strain and microstrain, and both of them will affect the diffracted peaks. The macro-strain will shift the diffraction peaks and affect the identification of specific peaks, while the micro-strain will affect the broadening of the peak and thus alter the size estimation of samples.

### 5.2.1 Macro-stain

Marco-strain means that the whole material is under directional residual tension or compression, so that the inter-planar spacing $d_{h k l}$ will be increase or decrease in certain direction. Therefore, the diffraction peak profile will be shifted [26].

Hammond [26] suggested that the macro-strain can be measured by rotating the specimen for various settings, and thus the direction and magnitude of the macro-strain can then be extrapolated.

### 5.2.2 Micro-strain

Micro-strain means that the directions or the magnitude of the internal strains are different from crystallites to crystallites [26]. These may be coming from, for instance faulting, dislocation [42], non-uniform lattice distortions [21], etc. The contribution to the broadening by micro-strain can be calculated by differentiate Bragg's Law (Equation 3.1) with $\lambda$ fixed [26]:

$$
\begin{equation*}
0=2 d \cos \theta \delta \theta+2 \sin \theta \delta d \tag{5.8}
\end{equation*}
$$

since micro-strain $\epsilon$ is equal to $\delta d / d$, so Equation 5.8 becomes [26]:

$$
\begin{equation*}
\epsilon=-\cot \theta \delta \theta \tag{5.9}
\end{equation*}
$$

Therefore the broadening due to micro-strain is:

$$
\begin{equation*}
B(2 \theta)=-\frac{2 \epsilon}{\cot \theta}=-2 \epsilon \tan \theta \tag{5.10}
\end{equation*}
$$

As mentioned by Hammond [26], broadening due to particle size varies as $1 \cos \theta$ while the broadening due to micro-strain varies as $1 \cot \theta$.

### 5.3 Thermal diffuse scattering

Atoms in the crystal vibrate, and these vibrations affect the diffracted intensities. The vibrations occurred for two reasons. The first one is purely quantum mechanical in origin from the uncertainty principle; the second one is due to elastic waves, or phonons, that are thermally excited in crystal. Equation 5.11 shows the scattered intensity with the consideration of these vibrations [43].
$I=\sum_{m} \sum_{n} f(\boldsymbol{Q}) e^{-M} e^{i \mathbf{Q} \cdot \boldsymbol{R}_{m}} f^{*}(\boldsymbol{Q}) e^{-M} e^{i \mathbf{Q} \cdot \boldsymbol{R}_{n}}+\sum_{m} \sum_{n} f(\boldsymbol{Q}) e^{-M} e^{i \boldsymbol{Q} \cdot \boldsymbol{R}_{m}} f^{*}(\boldsymbol{Q}) e^{-M} e^{i \boldsymbol{Q} \cdot \boldsymbol{R}_{n}}\left\{e^{Q^{2}\left(u_{Q m} u_{Q n}\right)}-1\right\}$
where $f$ is atomic form factor, $\mathbf{Q}$ is the wave-vector transfer, $e^{-M}$ is the Debye-Waller factor, $R_{n}$ is the time-averaged mean position for n'th atom, $u_{Q n}$ is the component of displacement parallel to $\boldsymbol{Q}$ for the n'th atom.

The first term in Equation 5.11 is the elastic Bragg scattering. The second term is thermal diffuse scattering (TDS) [43]. Figure 5.4 shows the idea of TDS contribution to the diffracted peak profile, but due to lack of experimental data, the TDS in Figure 5.4 is not to scale. The actual value of TDS can be found by experiment.

In general, TDS have to be eliminated to give a more accurate size estimation. Chapter 4 shows how background intensity was eliminated in this study.


Figure 5.4: TDS contribution to silver 200 reflection peak at $432^{\circ} \mathrm{C}$, where the TDS intensity is not to scale

## Chapter 6

## Growth kinetics

### 6.1 Isothermal kinetics

The Johnson-Mehl-Avrami (JMA) equation can be used to describe many types of isothermal solid state transformations [44]. The JMA equation has the form [44]:

$$
\begin{equation*}
x(t)=1-e^{-(K t)^{n}} \tag{6.1}
\end{equation*}
$$

where $x(t)$ is the volume fraction transformed after time $t, t$ is the elapsed time, $K$ is the reaction rate constant and $n$ is the kinetic exponent. By taking natural logarithm twice on both sides of Equation 6.1, JMA equation becomes:

$$
\begin{equation*}
\ln [-\ln (1-x(t))]=n \ln K-n \ln t \tag{6.2}
\end{equation*}
$$

In this study, $x(t)$ was assumed to be reached $100 \%$ when the estimated size of silver reached the plateau at $432^{\circ} \mathrm{C}$ and $x(t)$ from $380^{\circ} \mathrm{C}$ to $432^{\circ} \mathrm{C}$ was defined by Equation 6.3.

$$
\begin{equation*}
x(t)=\frac{D_{t}^{3}}{D_{100 \%}^{3}} \tag{6.3}
\end{equation*}
$$

where $D_{t}$ is the size estimated at elapsed time $t$ and $D_{100 \%}$ is the size estimated at $432^{\circ} \mathrm{C}$.
By plotting $\ln [-\ln (1-x(t))]$ against $\ln t$, the kinetics exponent n can be found by reading the slop of the plot and the reaction rate constant $K$ can be calculated from the intercept $n \ln (K)$.

However, the initial value of elapsed time $t$ at $380^{\circ} \mathrm{C}$ was unknown. The initial value of $t$ was assumed to be 0 at $380^{\circ} \mathrm{C}$ in the following part, so as to have a general picture of the growth kinetics.

Figure 6.1 and 6.2 show that, the plotting of $\ln [-\ln (1-x(t))]$ against $\ln t$ did not give a straight line for all five reflections by both FWHM method and integral breadth method. So, the growth kinetics could not be calculated directly from this model. However, it could be possible to assign two straight lines for each plots, one line for $\ln t<6.5$ and the other line for $\ln t<6.5$. Two sets of rate constants and kinetics exponents could possibly use to study the growth of silver.

One of the main reasons that plots did not show one straight line is that, the experiment was not done in non-isothermal condition, but readily increasing temperature at a rate of $2{ }^{\circ} \mathrm{C}$ per minutes. If this model will be used in future study, an isothermal experimental condition will be recommended.

### 6.2 Non-isothermal kinetics

According to Vazquez et al. [44], the non-isothermal kinetics model is much more complicated than the isothermal model, and contains some values that were not measured in this study. So, using the model in this study was not possible.


Figure 6.1: The plotting of $\ln [-\ln (1-x(t))]$ against $\ln t$ within the temperature range of $380^{\circ} \mathrm{C}$ to $432{ }^{\circ} \mathrm{C}$, for the five lowest reflections by defining peak breadth by FWHM


Figure 6.2: The plotting of $\ln [-\ln (1-x(t))]$ against $\ln t$ within the temperature range of $380^{\circ} \mathrm{C}$ to $432{ }^{\circ} \mathrm{C}$, for the five lowest reflections by defining peak breadth by integral breadth

### 6.3 Shape of silver crystallites

Although the growth kinetics cannot be found though the model, the shape of silver crystallites can be estimated by using the sizes calculated from Table 4.1. As the silver is FCC structure, ( hkl ) is equal to $[h k l]$. The estimated sizes of silver for each reflection can act as a probe for thickness in different directions. Therefore, the silver crystallites can be visualized. Since some of the directions cannot be measured due to extinction principle, the shapes of the silver crystallites were estimated by using thickness in [111], [200], [220] as the principal axes of ellipsoid. The estimated shape of silver crystallites by FWHM method was almost like a sphere shown in Figure 6.3, while the estimated shape of silver crystallites by integral breadth method was more like an ellipsoid shown in Figure 6.4.


Figure 6.3: Shape of silver crystallites by FWHM method, which was estimated by using thickness in [111], [200], [220] as the principal axes of ellipsoid


Figure 6.4: Shape of silver crystallites by integral breadth method, which was estimated by using thickness in [111], [200], [220] as the principal axes of ellipsoid

## Chapter 7

## Summary and suggestion to further

## investigation

### 7.1 Summary and conclusions

## Presence of silver in silver sulphate sample

In Chapter 3, Bragg's law and extinction rule were used to calculate the silver peak position as a function of temperature. Also, the effect of increasing temperature to silver lattice constants was also considered so as to have a more accurate estimation on the diffraction angle. Furthermore, exponential decay of synchrotron beams was also evaluated, so that the silver peak intensities could represent the true values.

The results in Chapter 3 show that, diffractograms in sample 1 and 3 did not show any silver peaks in the five lowest orders of reflections throughout the temperature range of this study. While sample 2 shows silver peaks in all five lowest orders of reflections. The difference between sample 2 and sample 1, 3 may be due to the difference resulted from the manufacturing process. The detailed mechanism concerning this point is not known yet and needs further investigation. Therefore, if silver size growth in silver sulphate during heating process is to be studied, sample 2 is recommended.

## Silver growth

Scherrer equation was applied in Chapter 4 to estimate the size of silver. The results show that silver crystallites were grown almost linearly with increasing temperature in all reflections within the temperature range of $380^{\circ} \mathrm{C}$ to $432^{\circ} \mathrm{C}$. The sizes of silver grew typically from around 50 $\pm 10 \mathrm{~nm}$ to $70 \pm 10 \mathrm{~nm}$ Then, silver crystallites stopped growing when temperature was over 432 ${ }^{\circ} \mathrm{C}$. Therefore, it is very likely that the phase transition of silver sulphate prohibits the growth of metallic silver rather than leads to the growth of silver, because the silver crystallites grew at temperature below the phase transition temperature and stopped at temperature near to the phase transition. The origin of the silver growth is suspected to be the disproportionation mechanism of silver sulphate.

## Comparison on FWHM and integral breadth

With reference to Chapter 4, both FWHM and integral breadth methods were tested in order to estimate the sizes of silver. The estimated sizes of silver by both methods grew at almost the same rate and pattern throughout the temperature range of the experiment. However, the estimated sizes have a constant difference throughout the temperature range of the experiment, and this might be due to the difference in defining peak breath and the values of Scherrer constant. For FWHM method, the sizes of silver grew typically from around $50 \pm 10 \mathrm{~nm}$ to $70 \pm 10$ nm . For integral breadth method, the sizes of silver grew typically from around $60 \pm 10 \mathrm{~nm}$ to 80 $\pm 10 \mathrm{~nm}$.

## Others factor leading to the broadening of diffraction peak

In this study, two of the main instrumental broadenings were coming from the finite pixel size of the detector and the energy dispersion of the synchrotron radiation. The uncertainties in FWHM due to finite pixel size of the detector were $0.04^{\circ} \pm 0.01^{\circ}$ to $0.10 \pm 0.01^{\circ}$.The magnitude of the instrumental broadening due to the energy dispersion of the synchrotron radiation was $10^{-5 \circ}$ and is insignificant. While other factors that lead to the broadening of diffraction peak like micro-strain were difficult to estimate quantitatively in this study, because of the lack of experimental data.

### 7.2 Suggestion to further investigation

## Origin of the silver metal

From Chapter 4, silver crystallites were found to be growing during the heating process of metallic silver. One of the suggestion to the growth of metallic silver during the heating process of silver sulphate is the disproportionation mechanism [45]:

$$
\begin{equation*}
2 \mathrm{Ag}_{2} \mathrm{SO}_{4} \rightarrow 2 \mathrm{Ag}+\mathrm{Ag}_{2} \mathrm{SO}_{4}+\mathrm{SO}_{2}+\mathrm{O}_{2} \tag{7.1}
\end{equation*}
$$

Disproportionation is a redox reaction that $\mathrm{Ag}_{2} \mathrm{SO}_{4}$ is both oxidized and reduced [46]. However, this is only a hypothesis right now, and there is no specific research that has been done on disproportionation of silver sulphate. Further investigation is needed to test this hypothesis.

## Reason to the stopping of silver growth

With reference to Chapter 4, silver growth stopped at around $432^{\circ} \mathrm{C}$ in all reflections. Phase transition of silver sulphate is suspected to cause the stop of silver growth, as the stopping temperature $\left(432^{\circ} \mathrm{C}\right)$ is very close to the phase transition temperature $\left(427^{\circ} \mathrm{C}\right)$.

To test the hypothesis, the silver sulphate powder have to be heated up from $380{ }^{\circ} \mathrm{C}$ to 530 ${ }^{\circ} \mathrm{C}$ and then cool down back to $380^{\circ} \mathrm{C}$. Then, it is needed to check if silver crystallites will grow again at temperature lower than the phase transition temperature of silver sulphate. If the silver crystallites grow again, then very likely the phase transition of sulphate silver is prohibiting the growth of metallic silver.

TEM can also be used for doing further investigation in the system, because TEM can be used to look at the local structure of the system, and might be able to find evidence to the site saturation of silver sulphate matrix during the silver growth.

## Appendix A

## Diffraction angles of silver

Table A. 1 and table A. 2 were calculated by Excel function, Data Table.

Table A.1: Diffraction angles of silver estimated by linear regression

|  | Reflections, $2 \theta$ (degree) |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Temperature, ${ }^{\circ} \mathrm{C}$ | 111 | 200 | 220 | 311 | 222 |  |
| 380 | 16.784 | 19.404 | 27.576 | 32.457 | 33.942 |  |
| 382 | 16.783 | 19.403 | 27.575 | 32.455 | 33.941 |  |
| 384 | 16.783 | 19.402 | 27.573 | 32.454 | 33.939 |  |
| 386 | 16.782 | 19.401 | 27.572 | 32.452 | 33.938 |  |
| 388 | 16.781 | 19.400 | 27.571 | 32.451 | 33.936 |  |
| 390 | 16.780 | 19.400 | 27.570 | 32.449 | 33.935 |  |
| 392 | 16.780 | 19.399 | 27.568 | 32.448 | 33.933 |  |
| 394 | 16.779 | 19.398 | 27.567 | 32.446 | 33.931 |  |
| 396 | 16.778 | 19.397 | 27.566 | 32.445 | 33.930 |  |
| 398 | 16.777 | 19.396 | 27.564 | 32.443 | 33.928 |  |
| 400 | 16.777 | 19.395 | 27.563 | 32.442 | 33.927 |  |
| 402 | 16.776 | 19.394 | 27.562 | 32.440 | 33.925 |  |
| 404 | 16.775 | 19.393 | 27.561 | 32.439 | 33.924 |  |
| 406 | 16.774 | 19.393 | 27.559 | 32.437 | 33.922 |  |
| 408 | 16.773 | 19.392 | 27.558 | 32.436 | 33.921 |  |
| 410 | 16.773 | 19.391 | 27.557 | 32.434 | 33.919 |  |
| 412 | 16.772 | 19.390 | 27.556 | 32.433 | 33.917 |  |
| 414 | 16.771 | 19.389 | 27.554 | 32.431 | 33.916 |  |
| 416 | 16.770 | 19.388 | 27.553 | 32.430 | 33.914 |  |
| 418 | 16.770 | 19.387 | 27.552 | 32.428 | 33.913 |  |
| 420 | 16.769 | 19.386 | 27.551 | 32.427 | 33.911 |  |
| 422 | 16.768 | 19.386 | 27.549 | 32.425 | 33.910 |  |
| 424 | 16.767 | 19.385 | 27.548 | 32.424 | 33.908 |  |
| 426 | 16.767 | 19.384 | 27.547 | 32.422 | 33.906 |  |
| 428 | 16.766 | 19.383 | 27.546 | 32.421 | 33.905 |  |
| 430 | 16.765 | 19.382 | 27.544 | 32.419 | 33.903 |  |
| 432 | 16.764 | 19.381 | 27.543 | 32.418 | 33.902 |  |
| 434 | 16.764 | 19.380 | 27.542 | 32.416 | 33.900 |  |
| 436 | 16.763 | 19.379 | 27.541 | 32.415 | 33.899 |  |
| 438 | 16.762 | 19.379 | 27.539 | 32.413 | 33.897 |  |
| 440 | 16.761 | 19.378 | 27.538 | 32.412 | 33.895 |  |
| 442 | 16.761 | 19.377 | 27.537 | 32.410 | 33.894 |  |
| 444 | 16.760 | 19.376 | 27.535 | 32.409 | 33.892 |  |
| 446 | 16.759 | 19.375 | 27.534 | 32.407 | 33.891 |  |
| 448 | 16.758 | 19.374 | 27.533 | 32.406 | 33.889 |  |
| 450 | 16.758 | 19.373 | 27.532 | 32.404 | 33.888 |  |
| 452 | 16.757 | 19.372 | 27.530 | 32.403 | 33.886 |  |
| 454 | 16.756 | 19.372 | 27.529 | 32.401 | 33.884 |  |
| 456 | 16.755 | 19.371 | 27.528 | 32.400 | 33.883 |  |
| 458 | 16.755 | 19.370 | 27.527 | 32.398 | 33.881 |  |
|  |  |  |  |  |  |  |


|  | Reflections, $2 \theta$ (degree) |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Temperature, ${ }^{\circ} \mathrm{C}$ | 111 | 200 | 220 | 311 | 222 |
| 460 | 16.754 | 19.369 | 27.525 | 32.397 | 33.880 |
| 462 | 16.753 | 19.368 | 27.524 | 32.395 | 33.878 |
| 464 | 16.752 | 19.367 | 27.523 | 32.394 | 33.877 |
| 466 | 16.752 | 19.366 | 27.522 | 32.392 | 33.875 |
| 468 | 16.751 | 19.365 | 27.520 | 32.391 | 33.874 |
| 470 | 16.750 | 19.364 | 27.519 | 32.389 | 33.872 |
| 472 | 16.749 | 19.364 | 27.518 | 32.388 | 33.870 |
| 474 | 16.749 | 19.363 | 27.517 | 32.386 | 33.869 |
| 476 | 16.748 | 19.362 | 27.515 | 32.385 | 33.867 |
| 478 | 16.747 | 19.361 | 27.514 | 32.383 | 33.866 |
| 480 | 16.746 | 19.360 | 27.513 | 32.382 | 33.864 |
| 482 | 16.745 | 19.359 | 27.512 | 32.380 | 33.863 |
| 484 | 16.745 | 19.358 | 27.510 | 32.379 | 33.861 |
| 486 | 16.744 | 19.357 | 27.509 | 32.377 | 33.859 |
| 488 | 16.743 | 19.357 | 27.508 | 32.376 | 33.858 |
| 490 | 16.742 | 19.356 | 27.507 | 32.374 | 33.856 |
| 492 | 16.742 | 19.355 | 27.505 | 32.373 | 33.855 |
| 494 | 16.741 | 19.354 | 27.504 | 32.371 | 33.853 |
| 496 | 16.740 | 19.353 | 27.503 | 32.370 | 33.852 |
| 498 | 16.739 | 19.352 | 27.502 | 32.368 | 33.850 |
| 500 | 16.739 | 19.351 | 27.500 | 32.367 | 33.849 |
| 502 | 16.738 | 19.350 | 27.499 | 32.365 | 33.847 |
| 504 | 16.737 | 19.350 | 27.498 | 32.364 | 33.845 |
| 506 | 16.736 | 19.349 | 27.497 | 32.362 | 33.844 |
| 508 | 16.736 | 19.348 | 27.495 | 32.361 | 33.842 |
| 510 | 16.735 | 19.347 | 27.494 | 32.360 | 33.841 |
| 512 | 16.734 | 19.346 | 27.493 | 32.358 | 33.839 |
| 514 | 16.733 | 19.345 | 27.491 | 32.357 | 33.838 |
| 516 | 16.733 | 19.344 | 27.490 | 32.355 | 33.836 |
| 518 | 16.732 | 19.343 | 27.489 | 32.354 | 33.835 |
| 520 | 16.731 | 19.343 | 27.488 | 32.352 | 33.833 |
| 522 | 16.730 | 19.342 | 27.486 | 32.351 | 33.831 |
| 524 | 16.730 | 19.341 | 27.485 | 32.349 | 33.830 |
| 526 | 16.729 | 19.340 | 27.484 | 32.348 | 33.828 |
| 528 | 16.728 | 19.339 | 27.483 | 32.346 | 33.827 |
| 530 | 16.727 | 19.338 | 27.481 | 32.345 | 33.825 |
| 532 | 16.727 | 19.337 | 27.480 | 32.343 | 33.824 |
| 534 | 16.726 | 19.337 | 27.479 | 32.342 | 33.822 |
| 536 | 16.725 | 19.336 | 27.478 | 32.340 | 33.820 |
| 538 | 16.724 | 19.335 | 27.476 | 32.339 | 33.819 |
|  | 16.724 | 19.334 | 27.475 | 32.337 | 33.817 |
|  |  |  |  |  |  |

Table A.2: Diffraction angles of silver estimated by exponential regression

|  | Reflections, $2 \theta$ (degree) |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Temperature, ${ }^{\circ} \mathrm{C}$ | 111 | 200 | 220 | 311 | 222 |
| 380 | 16.784 | 19.404 | 27.576 | 32.457 | 33.943 |
| 382 | 16.784 | 19.404 | 27.575 | 32.456 | 33.942 |
| 384 | 16.783 | 19.403 | 27.574 | 32.454 | 33.940 |
| 386 | 16.782 | 19.402 | 27.573 | 32.453 | 33.939 |
| 388 | 16.781 | 19.401 | 27.571 | 32.451 | 33.937 |
| 390 | 16.781 | 19.400 | 27.570 | 32.450 | 33.935 |
| 392 | 16.780 | 19.399 | 27.569 | 32.448 | 33.934 |
| 394 | 16.779 | 19.398 | 27.568 | 32.447 | 33.932 |
| 396 | 16.778 | 19.397 | 27.566 | 32.445 | 33.931 |
| 398 | 16.778 | 19.397 | 27.565 | 32.444 | 33.929 |
| 400 | 16.777 | 19.396 | 27.564 | 32.442 | 33.928 |
| 402 | 16.776 | 19.395 | 27.563 | 32.441 | 33.926 |
| 404 | 16.775 | 19.394 | 27.561 | 32.439 | 33.924 |
| 406 | 16.775 | 19.393 | 27.560 | 32.438 | 33.923 |
| 408 | 16.774 | 19.392 | 27.559 | 32.436 | 33.921 |
| 410 | 16.773 | 19.391 | 27.558 | 32.435 | 33.920 |
| 412 | 16.772 | 19.390 | 27.556 | 32.433 | 33.918 |
| 414 | 16.772 | 19.389 | 27.555 | 32.432 | 33.917 |
| 416 | 16.771 | 19.389 | 27.554 | 32.430 | 33.915 |
| 418 | 16.770 | 19.388 | 27.552 | 32.429 | 33.913 |
| 420 | 16.769 | 19.387 | 27.551 | 32.427 | 33.912 |
| 422 | 16.769 | 19.386 | 27.550 | 32.426 | 33.910 |
| 424 | 16.768 | 19.385 | 27.549 | 32.424 | 33.909 |
| 426 | 16.767 | 19.384 | 27.547 | 32.423 | 33.907 |
| 428 | 16.766 | 19.383 | 27.546 | 32.421 | 33.906 |
| 430 | 16.766 | 19.382 | 27.545 | 32.420 | 33.904 |
| 432 | 16.765 | 19.382 | 27.544 | 32.418 | 33.903 |
| 434 | 16.764 | 19.381 | 27.542 | 32.417 | 33.901 |
| 436 | 16.763 | 19.380 | 27.541 | 32.415 | 33.899 |
| 438 | 16.763 | 19.379 | 27.540 | 32.414 | 33.898 |
| 440 | 16.762 | 19.378 | 27.539 | 32.412 | 33.896 |
| 442 | 16.761 | 19.377 | 27.537 | 32.411 | 33.895 |
| 444 | 16.760 | 19.376 | 27.536 | 32.409 | 33.893 |
| 446 | 16.759 | 19.375 | 27.535 | 32.408 | 33.892 |
| 448 | 16.759 | 19.375 | 27.534 | 32.406 | 33.890 |
| 450 | 16.758 | 19.374 | 27.532 | 32.405 | 33.888 |
| 452 | 16.757 | 19.373 | 27.531 | 32.404 | 33.887 |
| 454 | 16.756 | 19.372 | 27.530 | 32.402 | 33.885 |
| 456 | 16.756 | 19.371 | 27.529 | 32.401 | 33.884 |
| 458 | 16.755 | 19.370 | 27.527 | 32.399 | 33.882 |
|  |  |  |  |  |  |
|  |  |  |  |  |  |
|  |  |  |  |  |  |


| Temperature, ${ }^{\circ} \mathrm{C}$ | Reflections, $2 \theta$ (degree) |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | 111 | 200 | 220 | 311 | 222 |
| 460 | 16.754 | 19.369 | 27.526 | 32.398 | 33.881 |
| 462 | 16.753 | 19.368 | 27.525 | 32.396 | 33.879 |
| 464 | 16.753 | 19.368 | 27.524 | 32.395 | 33.877 |
| 466 | 16.752 | 19.367 | 27.522 | 32.393 | 33.876 |
| 468 | 16.751 | 19.366 | 27.521 | 32.392 | 33.874 |
| 470 | 16.750 | 19.365 | 27.520 | 32.390 | 33.873 |
| 472 | 16.750 | 19.364 | 27.519 | 32.389 | 33.871 |
| 474 | 16.749 | 19.363 | 27.517 | 32.387 | 33.870 |
| 476 | 16.748 | 19.362 | 27.516 | 32.386 | 33.868 |
| 478 | 16.747 | 19.361 | 27.515 | 32.384 | 33.867 |
| 480 | 16.747 | 19.361 | 27.513 | 32.383 | 33.865 |
| 482 | 16.746 | 19.360 | 27.512 | 32.381 | 33.863 |
| 484 | 16.745 | 19.359 | 27.511 | 32.380 | 33.862 |
| 486 | 16.744 | 19.358 | 27.510 | 32.378 | 33.860 |
| 488 | 16.744 | 19.357 | 27.508 | 32.377 | 33.859 |
| 490 | 16.743 | 19.356 | 27.507 | 32.375 | 33.857 |
| 492 | 16.742 | 19.355 | 27.506 | 32.374 | 33.856 |
| 494 | 16.741 | 19.354 | 27.505 | 32.372 | 33.854 |
| 496 | 16.741 | 19.354 | 27.503 | 32.371 | 33.852 |
| 498 | 16.740 | 19.353 | 27.502 | 32.369 | 33.851 |
| 500 | 16.739 | 19.352 | 27.501 | 32.368 | 33.849 |
| 502 | 16.738 | 19.351 | 27.500 | 32.366 | 33.848 |
| 504 | 16.738 | 19.350 | 27.498 | 32.365 | 33.846 |
| 506 | 16.737 | 19.349 | 27.497 | 32.363 | 33.845 |
| 508 | 16.736 | 19.348 | 27.496 | 32.362 | 33.843 |
| 510 | 16.735 | 19.347 | 27.495 | 32.360 | 33.842 |
| 512 | 16.735 | 19.347 | 27.493 | 32.359 | 33.840 |
| 514 | 16.734 | 19.346 | 27.492 | 32.357 | 33.838 |
| 516 | 16.733 | 19.345 | 27.491 | 32.356 | 33.837 |
| 518 | 16.732 | 19.344 | 27.490 | 32.354 | 33.835 |
| 520 | 16.732 | 19.343 | 27.488 | 32.353 | 33.834 |
| 522 | 16.731 | 19.342 | 27.487 | 32.351 | 33.832 |
| 524 | 16.730 | 19.341 | 27.486 | 32.350 | 33.831 |
| 526 | 16.729 | 19.340 | 27.485 | 32.348 | 33.829 |
| 528 | 16.729 | 19.340 | 27.483 | 32.347 | 33.827 |
| 530 | 16.728 | 19.339 | 27.482 | 32.345 | 33.826 |
| 532 | 16.727 | 19.338 | 27.481 | 32.344 | 33.824 |
| 534 | 16.726 | 19.337 | 27.480 | 32.342 | 33.823 |
| 536 | 16.725 | 19.336 | 27.478 | 32.341 | 33.821 |
| 538 | 16.725 | 19.335 | 27.477 | 32.339 | 33.820 |
| 540 | 16.724 | 19.334 | 27.476 | 32.338 | 33.818 |

## Appendix B

## Scaling factors of exponential decay of synchrotron beam

Table B. 1 was calculated by Equation 3.6 with half life equals to 10 hours and the elapsed time before the first recorded data equals to 1500 s and time interval between data sets equals to 60 s.

Table B.1: Scaling factor $K$ for the samples at different elapsed time

| Data set | Elapsed time, s | Scaling factor $K$ | Data set | Elapsed time, s | Scaling factor $K$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1500 | 1.0425 | 16 | 2400 | 1.0689 |
| 2 | 1560 | 1.0443 | 17 | 2460 | 1.0707 |
| 3 | 1620 | 1.0460 | 18 | 2520 | 1.0725 |
| 4 | 1680 | 1.0478 | 19 | 2580 | 1.0743 |
| 5 | 1740 | 1.0495 | 20 | 2640 | 1.0761 |
| 6 | 1800 | 1.0513 | 21 | 2700 | 1.0779 |
| 7 | 1860 | 1.0530 | 22 | 2760 | 1.0797 |
| 8 | 1920 | 1.0548 | 23 | 2820 | 1.0815 |
| 9 | 1980 | 1.0565 | 24 | 2880 | 1.0833 |
| 10 | 2040 | 1.0583 | 25 | 2940 | 1.0851 |
| 11 | 2100 | 1.0601 | 26 | 3000 | 1.0869 |
| 12 | 2160 | 1.0618 | 27 | 3060 | 1.0887 |
| 13 | 2220 | 1.0636 | 28 | 3120 | 1.0905 |
| 14 | 2280 | 1.0654 | 29 | 3180 | 1.0924 |
| 15 | 2340 | 1.0672 | 30 | 3240 | 1.0942 |


| Data set | Elapsed time, s | Scaling factor $K$ | Data set | Elapsed time, s | Scaling factor $K$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 31 | 3300 | 1.0960 | 65 | 5340 | 1.1599 |
| 32 | 3360 | 1.0978 | 66 | 5400 | 1.1618 |
| 33 | 3420 | 1.0997 | 67 | 5460 | 1.1638 |
| 34 | 3480 | 1.1015 | 68 | 5520 | 1.1657 |
| 35 | 3540 | 1.1033 | 69 | 5580 | 1.1677 |
| 36 | 3600 | 1.1052 | 70 | 5640 | 1.1696 |
| 37 | 3660 | 1.1070 | 71 | 5700 | 1.1716 |
| 38 | 3720 | 1.1089 | 72 | 5760 | 1.1735 |
| 39 | 3780 | 1.1107 | 73 | 5820 | 1.1755 |
| 40 | 3840 | 1.1126 | 74 | 5880 | 1.1774 |
| 41 | 3900 | 1.1144 | 75 | 5940 | 1.1794 |
| 42 | 3960 | 1.1163 | 76 | 6000 | 1.1814 |
| 43 | 4020 | 1.1181 | 77 | 6060 | 1.1833 |
| 44 | 4080 | 1.1200 | 78 | 6120 | 1.1853 |
| 45 | 4140 | 1.1219 | 79 | 6180 | 1.1873 |
| 46 | 4200 | 1.1237 | 80 | 6240 | 1.1893 |
| 47 | 4260 | 1.1256 | 81 | 6300 | 1.1912 |
| 48 | 4320 | 1.1275 | 82 | 6360 | 1.1932 |
| 49 | 4380 | 1.1294 | 83 | 6420 | 1.1952 |
| 50 | 4440 | 1.1313 | 84 | 6480 | 1.1972 |
| 51 | 4500 | 1.1331 | 85 | 6540 | 1.1992 |
| 52 | 4560 | 1.1350 | 86 | 6600 | 1.2012 |
| 53 | 4620 | 1.1369 | 87 | 6660 | 1.2032 |
| 54 | 4680 | 1.1388 | 88 | 6720 | 1.2052 |
| 55 | 4740 | 1.1407 | 89 | 6780 | 1.2072 |
| 56 | 4800 | 1.1426 | 90 | 6840 | 1.2092 |
| 57 | 4860 | 1.1445 | 91 | 6900 | 1.2113 |
| 58 | 4920 | 1.1464 | 92 | 6960 | 1.2133 |
| 59 | 4980 | 1.1484 | 93 | 7020 | 1.2153 |
| 60 | 5040 | 1.1503 | 94 | 7080 | 1.2173 |
| 61 | 5100 | 1.1522 | 95 | 7140 | 1.2194 |
| 62 | 5160 | 1.1541 | 96 | 7200 | 1.2214 |
| 63 | 5220 | 1.1560 | 97 | 7260 | 1.2234 |
| 64 | 5280 | 1.1580 | 98 | 7320 | 1.2255 |
|  |  |  |  |  |  |
|  |  |  |  |  |  |
|  |  |  |  |  |  |

## Appendix C

## Maximum instrumental broadening due to

 energy dispersionTable C.1: Instrumental broadening due to energy dispersion of synchrotron radiation

|  | Reflections, $2 \theta$ (degree) |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Temperature | 111 | 200 | 220 | 311 | 222 |
| 380 | $2.89 \times 10^{-5}$ | $3.33 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 382 | $2.89 \times 10^{-5}$ | $3.33 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 384 | $2.89 \times 10^{-5}$ | $3.33 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 386 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 388 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 390 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 392 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 394 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 396 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 398 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 400 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 402 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 404 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 406 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 408 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 410 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 412 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.59 \times 10^{-5}$ |
| 414 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 416 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 418 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 420 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 422 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 424 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 426 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 428 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 430 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 432 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 434 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 436 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 438 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 440 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 442 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.37 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 444 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 446 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 448 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 450 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 452 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 454 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 456 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 458 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |


|  | Reflections, $2 \theta$ (degree) |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Temperature | 111 | 200 | 220 | 311 | 222 |
| 460 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 462 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 464 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 466 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 468 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 470 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 472 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 474 | $2.89 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.63 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 476 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 478 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 480 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 482 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 484 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 486 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 488 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 490 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 492 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 494 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 496 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 498 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 500 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 502 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 504 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 506 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 508 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 510 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.58 \times 10^{-5}$ |
| 512 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 514 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 516 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 518 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 520 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 522 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 524 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 526 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 528 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 530 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 532 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 534 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 536 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 538 | $2.88 \times 10^{-5}$ | $3.32 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |
| 540 | $2.88 \times 10^{-5}$ | $3.31 \times 10^{-5}$ | $4.62 \times 10^{-5}$ | $5.36 \times 10^{-5}$ | $5.57 \times 10^{-5}$ |

## Appendix D

## Mathematica Code

## Listing D.1: Import data set

```
ClearAII["Global'*"]
testlist =sample=re111=re200=re220=re311=re222=re111IntenVsTemp=re200IntenVsTemp=
        re22OIntenVsTemp=re311IntenVsTemp=re222IntenVsTemp=Table[1,{i,3}];
(* testlist [[1]] to testlist [[3]] represent sample 1, sample 2 and sample 3;
teslist [[ALI ]][[1]] to [[83]] and then to [[98]] respresent temperature from 380\[Degree]C to 540\[
        Degree]C and then back to 510\[Degree\C;
for each teslist, it contains three values: 1. 2\[Theta] in degrees 2. recorded diffracted intensity (
        arbitrary units) 3. ucertainty of the intensity measurement *)
testlist [[1]]= Table[Import["C:\\Users\\roy\\Desktop\\thesis\\Ag_2SO_4_powder\\ramp_1\\dat\\
        Agramp1s"<>ToString[i]<>".dat","Data"],{i,98}];
testlist [[2]]= Table[Import["C:\\Users\\roy\\Desktop\\thesis\\Ag_2SO_4_powder\\ramp_2\\dat\\
        Agramp2s"<>ToString[i]<>".dat","Data"],{i,98}];
testlist [[3]]= Table[Import["C:\\Users\\roy\\\esktop\\thesis\\Ag_2SO_4_powder\\ramp_3\\dat\\
        Agramp3s"<>ToString[i]<>".dat","Data"],{i,98}];
(*Scaling of the synchrotron beams because of the reduction in intensity of synchrotron beams along time
        *)
scaling =Flatten[Import["C:\\Users\\roy\\Desktop\\thesis\\scaling factor result . xlsx " ]];
(*Extract the values of the recorded diffracted intensity (arbitrary units) and the corresponding
    diffraction angle (2\[Theta])*)
```

```
Do[sample[[n]]=Table[Thread[{testlist [[n ]][[ i ]][[ All ,1]], testlist [[n ]][[ i ]][[ All ,2]]* scaling [[ i ]]}],{ i
        ,98}],{ n ,3}];
(*Data list of peak intensity and the corresponding diffraction angle (2\[Theta]) for 5 lowest
    reflection of silver of the three different samples*)
Do[re111[[n]]=Table[Flatten[MaximalBy[Select[sample[[n]][[ i]],(16.65<#[[1]]<16.95)&],Last ]],{ i ,98}];,{ n
        ,3}];
Do[re200[[n]]=Table[Flatten[MaximalBy[Select[sample[[n]][[i]],(19.25<#[[1]]<19.55)&],Last ]],{ i ,98}];,{ n
        ,3}];
Do[re220[[n]]=Table[Flatten[MaximalBy[Select[sample[[n]][[ i]],(27.35<#[[1]]<27.75)&],Last ]],{ i ,98}];,{ n
        ,3}];
Do[re311[[n]]=Table[Flatten[MaximalBy[Select[sample[[n]][[i]],(32.25<#[[1]]<32.65)&],Last ]],{ i ,98}];,{ n
        ,3}];
Do[re222[[n]]=Table[Flatten[MaximalBy[Select[sample[[n]][[ i]],(33.75<#[[1]]<34.05)&],Last ]],{ i ,98}];,{ n
        ,3}];
(* fix side peak*)
re220 [[2,1;;15]]= Table[Flatten[MaximalBy[Select[sample [[2]][[ i]],(27.55<#[[1]]<27.75)&],Last ]],{ i ,15}];
(*Data list of the temperature*)
Temp=Flatten[Import["C:\\Users\\roy\\\Desktop\\thesis\\temp.xlsx"]];
(*Data list of time in minutes*)
time=Flatten[Import["C:\\Users\\roy\\\Desktop\\thesis\\timeintervaloframpminutes.xlsx"]];
(*Reflections intensity vs Temperature*)
Do[re111IntenVsTemp[[n]]=Thread[{Temp,re111[[n]][[All ,2]]}],{n,3}];
Do[re200IntenVsTemp[[n]]=Thread[{Temp,re200[[n]][[All ,2]]}],{n,3}];
Do[re220IntenVsTemp[[n]]=Thread[{Temp,re220[[n]][[All ,2]]}],{n,3}];
Do[re311IntenVsTemp[[n]]=Thread[{Temp,re311[[n]][[All ,2]]}],{n,3}];
Do[re222IntenVsTemp[[n]]=Thread[{Temp,re222[[n]][[All ,2]]}],{ n ,3}];
```

Listing D.2: Chapter 2

(*Example of diffractograms of sample2 at 380\[Degree]C*)
ListPlot[\{sample [[2,1]]\}, PlotRange->All,Joined-> True,AxesLabel->\{"2\[Theta](degree)","Intensity ( arbitrary units)"\},ImageSize->500,AspectRatio->0.5]

## Listing D.3: Chapter 3

(*Plot: Subscript [Ag, 2]Subscript [SO, 4] peak nearby silver 220 reflection peak*)

ListPlot[sample [[2,1;;4]], PlotRange $->\{\{27.35,27.75\},\{20,80\}\}$, Joined $->$ True, Mesh $->$ All,ImageSize ->500,PlotMarkers->Automatic,PlotLegends->\{"380\[Degree]C","382\[Degree]C","384\[Degree]C"," $386 \backslash[$ Degree]C"\},PlotLegends->\{"380\[Degree]C","382\[Degree]C","384\[Degree]C","386\[Degree]C"\}, AxesLabel $->\{$ " $2 \backslash[$ Theta $]($ degree)", "Intensity (arbitrary units)" $\}$,Epilog $->\{$ Text["Subscript[Ag, 2] Subscript[SO, 4] peak",Scaled[\{.52,.94\}]],Text["Silver peak",Scaled[\{.72,.81\}]]\}]
(*Plot: To show the influence without proper scaling of the silver 111 reflection intensity *) sample2withoutScaling=Table[Thread[\{testlist [[2]]][[ i ]][[ All ,1]], testlist [[2]][[ i ]][[ All ,2]]\}],\{ i,98\}]; re111withoutScaling=Table[Flatten[MaximalBy[Select[sample2withoutScaling[[i]],(16.65<\#[[1]]<16.95)\&], Last]],\{ i ,98\}];
re111withoutScalingIntenVsTemp=Thread[\{Temp,re111withoutScaling[[All,2]]\}];ListPlot[\{ re111withoutScalingIntenVsTemp,re111IntenVsTemp[[2]]\},PlotRange->All,Mesh->All,PlotMarkers->\{ Automatic, 7$\}$, Joined-> True,AxesLabel-> "Temperature( $\backslash[$ Degree]C)","Intensity (arbitrary units)"\}, PlotLegends-> \{"Sample2 without scaling","Sample2 with scaling"\},ImageSize->500]

ListPlot [\{re111withoutScalingIntenVsTemp,re111IntenVsTemp[[2]]\},PlotRange->\{\{420,540\},\{500,700\}\}, Joined->True,Mesh->All,PlotMarkers->\{Automatic, 7$\}$,AxesLabel->\{"Temperature( $\backslash[$ Degree]C)"," Intensity (arbitrary units)"\},PlotLegends->\{"Sample2 without scaling", "Sample2 with scaling"\}, ImageSize->500]

(*Plot: : Intensity of silver 111 reflection of the 3 different samples along with temperature*) ListPlot[re111IntenVsTemp,PlotRange $->$ All,Joined $->$ True,Mesh $->$ All,PlotMarkers $->$ Automatic, 7 \}, AxesLabel->\{"Temperature(\[Degree]C)","Intensity (arbitrary units)"\},PlotLegends->\{"Sample1"," Sample2","Sample3"\},ImageSize->500]

(*Plot: Diffractogram to examining the presence of silver in the samples at temperature530\[PlusMinus] Cin 111 reflection *)

ListPlot[sample[[All ,83]], PlotRange $->\{\{16.60,16.85\},\{0,700\}\}$, Joined $->$ True,Mesh $->$ All,PlotMarkers $->$ Automatic,AxesLabel->\{"2\[Theta](degree)","Intensity (arbitrary units)"\},PlotLegends->\{"Sample1"," Sample2", "Sample3"\},ImageSize->500,PlotMarkers->Automatic]

Listing D.4: Chapter 4

```
(*Showing the diffraction peak broadening by crystallites size*)
Plot[ Sin[500x]^2/Sin[x]^2,{x,0,3Pi/2},Ticks-> {{0,Pi/2,Pi,3 Pi/2},Automatic},PlotRange->AlI,
    PlotLegends->"N = 500"]
```

Plot $\left[\operatorname{Sin}[5 x]^{\wedge} 2 / \operatorname{Sin}[x]^{\wedge} 2,\{x, 0,3 \mathbf{P i} / 2\}\right.$, Ticks $->\{\{0, \mathbf{P i} / 2, \mathbf{P i}, 3 \mathbf{P i} / 2\}$, Automatic $\}$, PlotRange $->$ All,
PlotLegends->"N = 5"]
(*Size estimation by FWHM*)
data1=Table[Select[sample [[2]][[ i ]],(16.6<\#[[1]]<16.85)\&], \{ i ,98\}];
$\operatorname{For}[\mathrm{i}=1, \mathrm{i}<99, \mathrm{i}++, \operatorname{data} 1[[i]]=\operatorname{Thread}[\{\operatorname{data} 1[[\mathrm{i}]][[$ All ,1]], data1[[i ]][[ All ,2]]- Mean[data1[[i ]][[1;;3,2]]]\}]];
nlm1 $=$ Table[1,\{i,98\}];
nlm1 $=$ Table[NonlinearModelFit[data1[[i ]], A $* \operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[\mathrm{Mu}])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma $], 0.05\},\{\backslash[\mathrm{Mu}], 16.75\},\{\mathrm{A}, 600\}\}, \times],\{i, 98\}] ;$

FWHM1=Table[nlm1[[i]]["ParameterTableEntries"][[1,1]]*2*Sqrt[2*Log[2]],\{i,98\}];
data2=Table[Select[sample [[2]][[ i ]],(19.25<\#[[1]]<19.55)\&], \{i ,98\}];
For $[i=1, i<99, i++$, data2[[i]]]=Thread[\{data2[[i ]][[ All ,1]], data2[[i $]][[$ All ,2]]- Mean[data2[[i ]][[-3;,,2]]]\}]];
nlm2 $=$ Table[1, $\{\mathrm{i}, 98\}$ ];
nlm2 $=$ Table[NonlinearModelFit[data2[[i]], A $* \operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[M u])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma
$], 0.05\},\{\backslash[\mathrm{Mu}], 19.4\},\{\mathrm{A}, 600\}\}, \mathrm{x}],\{i, 98\}] ;$
FWHM2=Table[nIm2[[i]]["ParameterTableEntries"][[1,1]]*2*Sqrt[2*Log[2]],\{i ,98\}];
data3 $=$ Table $[\{1,1\},\{i, 98\}] ;$
data3 $[[1 ; ; 11]]=$ Table[Select[sample [[2]]][[ i]],(27.56<\#[[1]]<27.75) \&], \{i ,1,11\}]; data3 [[12;;22]] = Table[
Select[sample [[2]][[ i]],(27.55<\#[[1]]<27.75)\&], \{i ,12,22\}]; data3 [[23;;30]]= Table[Select[sample [[2]][[ i ]],(27.54<\#[[1]]<27.75)\&],\{i ,23,30\}]; data3 [[31;;98]]= Table[Select[sample [[2]]][[ i
]],(27.35<\#[[1]]<27.75)\&], \{i ,31,98\}];
For $[\mathrm{i}=1, \mathrm{i}<99, \mathrm{i}++$, data3[[i]]=Thread[\{data3[[i ]][[ All ,1]], data3[[i ]][[ All ,2]]- Mean[data3[[i
]][[-3;;,2]]]\}];];
nlm3 $=$ Table[1,\{i,98\}];
nlm3 [[1;;11]] = Table[NonlinearModelFit[data3 [[i ] ], A * Exp[(-1) $\left.*(1 / 2) *(x-\backslash[\mathrm{Mu}])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma], 0.05\}, $\{\backslash[\mathrm{Mu}], 27.55\},\{\mathrm{A}, 200\}\}, \mathrm{x}],\{i, 1,11\}] ; n \operatorname{lm} 3[[11 ; ; 22]]=$ Table[NonlinearModelFit[data3[[i ]], A * $\operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[\mathrm{Mu}])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma $\left.\left.], 0.05\},\{\backslash[\mathrm{Mu}], 27.55\},\{\mathrm{A}, 200\}\}, \mathrm{x}\right],\{i, 11,22\}\right] ;$ nlm3 [[23;;30]] = Table[NonlinearModelFit[data3[[i ]], A $* \operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[\mathrm{Mu}])^{\wedge} 2 *(1 / \backslash[\right.$ Sigma $])$ ^2], $\{\{\backslash[$ Sigma $], 0.05\},\{\backslash[\mathrm{Mu}], 27.55\},\{\mathrm{A}, 200\}\}, x],\{i, 23,30\}] ; n \operatorname{lm} 3[[31 ; ; 98]]=$ Table[NonlinearModelFit[data3 $[[\mathrm{i}]], \mathrm{A} * \operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[\mathrm{Mu}])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma $\left.], 0.05\},\{\backslash[\mathrm{Mu}], 27.55\},\{\mathrm{A}, 200\}\}, \mathrm{x}\right],\{\mathrm{i}$ ,31,98\}];
FWHM3=Table[nlm3[[i]]]["ParameterTableEntries"][[1,1]]*2*Sqrt[2*Log[2]],\{i ,98\}];
data4=Table[Select[sample [[2]][[ i]],(32.25<\#[[1]]<32.65)\&], \{i ,98\}]
For $[i=1, i<99, i++, \operatorname{data} 4[[i]]=\operatorname{Thread}[\{\operatorname{data} 4[[i]][[$ All ,1]], data4[[ $i]][[$ All ,2]]- Mean[data4[[i
]][[-3;;,2]]]\}]];
nlm4 $=$ Table[1,\{i,98\}];
nlm4 $=$ Table[NonlinearModelFit[data4[[i ]], A $* \operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[M u])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma
$], 0.05\},\{\backslash[\mathrm{Mu}], 32.4\},\{\mathrm{A}, 200\}\}, \mathrm{x}],\{i, 98\}] ;$
FWHM4 = Table[nIm4[[i]]["ParameterTableEntries"][[1,1]]*2*Sqrt[2*Log[2]],\{i ,98\}];
data5=Table[Select[sample [[2]][[ i ]],(33.75<\#[[1]]<34.15)\&],\{i ,98\}];
For $[i=1, i<99, i++, \operatorname{data5}[[i]]=\operatorname{Thread}[\{\operatorname{data5}[[i]][[$ All ,1]], data5[[i ]][[ All ,2]]- Mean[data5[[i
]][[-3;;,2]]]\}]];
nlm5 $=$ Table[1, $\{i, 98\}]$;
$n \operatorname{lm} 5=$ Table[NonlinearModelFit[data5[[i]], A $* \operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[\mathrm{Mu}])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma
$], 0.05\},\{\backslash[\mathrm{Mu}], 33.95\},\{\mathrm{A}, 12\}\}, x],\{i, 98\}] ;$
FWHM5 = Table[nlm5[[i]]["ParameterTableEntries"][[1,1]]*2*Sqrt[2*Log[2]],\{i,98\}];
(*y-axis: thickness*)
thickness $1=\left(0.8551 *\left(0.6941 / 10^{\wedge} 10\right)\right) /((F W H M 1 * P i / 180) * \operatorname{Cos}[r e 111[[2$, All,1] $] / 2$ Degree $])$;
thickness2 $=\left(0.8859 *\left(0.6941 / 10^{\wedge} 10\right)\right) /((\mathrm{FWHM} 2 * \mathrm{Pi} / 180) * \operatorname{Cos}[r e 200[[2$, All,1]]/2Degree] $) ;$
thickness $3=\left(0.8340 *\left(0.6941 / 10^{\wedge} 10\right)\right) /((F W H M 3 * P i / 180) * \operatorname{Cos}[r e 220[[2$, All,1] $] / 2$ Degree $]) ;$
thickness $4=\left(0.9082 *\left(0.6941 / 10^{\wedge} 10\right)\right) /((\mathrm{FWHM} 4 * \mathrm{Pi} / 180) * \operatorname{Cos}[$ re311[[2,All,1]]/2Degree $])$;
thickness5 $=\left(0.8551 *\left(0.6941 / 10^{\wedge} 10\right)\right) /((F W H M 5 * P i / 180) * \operatorname{Cos}[r e 222[[2, A l l, 1]] / 2$ Degree $]) ;$
(*thickness vs temperature, assuming cubic crystallite and SC with Langfor's table*)
thickness1VsTemp=Thread[\{Temp,thickness1\}];
thickness2VsTemp=Thread[\{Temp,thickness2\}];
thickness3VsTemp=Thread[\{Temp,thickness3\}];
thickness4VsTemp=Thread[\{Temp,thickness4\}];
thickness5VsTemp=Thread[\{Temp,thickness5\}];
(*thickness vs time, assuming cubic crystallite and SC with Langfor's table*)
thickness1VsTime=Thread[\{time,thickness1*10^9\}];
thickness2VsTime=Thread[\{time,thickness2*10^9\}];
thickness3VsTime=Thread[\{time,thickness3*10^9\}];
thickness4VsTime=Thread[\{time,thickness $\left.4 * 10^{\wedge} 9\right\}$ ];

```
thickness5VsTime=Thread[{time,thickness5*10^9}];
(*thickness vs time, assuming spherical crystallite and SC 0.94 *)
thickness1CS = (0.94*(0.6941/10^10))/((FWHM1*Pi/180)*Cos[re111[[2,All,1]]/2Degree]);
thickness2CS =(0.94*(0.6941/10^10))/((FWHM2*Pi/180)*Cos[re200[[2,All,1]]/2Degree]);
thickness3CS =(0.94*(0.6941/10^10))}/((\textrm{FWHM}3*Pi/180)*\operatorname{Cos[re220[[2,All,1]]/2Degree]})
thickness4CS =(0.94*(0.6941/10^10))/((FWHM4*Pi/180)*Cos[re311[[2,All,1]]/2Degree]);
thickness5CS =(0.94*(0.6941/10^10))/((FWHM5*Pi/180)*Cos[re222[[2,All,1]]/2Degree]);
```

thickness1CSVsTime=Thread[\{time,thickness1CS*10^9\}];
thickness2CSVsTime=Thread[\{time,thickness2CS*10^9\}];
thickness3CSVsTime=Thread[\{time,thickness3CS*10^9\}];
thickness4CSVsTime=Thread[\{time,thickness4CS*10^9\}];
thickness5CSVsTime=Thread[\{time,thickness5CS*10^9\}];

(*To show the Gaussian fitting, e.g. 200 reflection at 398\[Degree]C*)
ListPlot[data2 [[10]], PlotStyle $->$ Red,PlotRange $->$ All, AxesLabel $->\{$ " $2 \backslash[$ Theta $]($ degree)", "Intensity (
arbitrary units)" $\}$,ImageSize $->500$ ]
Show[ListPlot[Thread[\{data2[[10]][[All ,1]], data2 [[10]][[ All ,2]]-Mean[data2 [[10]][[1;;3,2]]]\}], PlotStyle
$->$ Red, PlotRange $->$ All, ImageSize $->500$, AxesLabel $->\{$ " $2 \backslash$ [Theta](degree)", "Intensity (arbitrary units
)" $\}]$, Plot[nlm2[[10]][x],\{x,19.25,19.55\},PlotRange->All,ImageSize->500,AxesLabel->\{" $2 \backslash[$ Theta](degree)","Intensity (arbitrary units)"\}]]
(*Show cancelling of background intensity *)
databg2=Table[Select[sample[[2]][[ i]],(19.25<\#[[1]]<19.55)\&], \{i ,98\}];
For[ $i=1, i<99, i++$, databg2[[i]] $=$ Thread[\{databg2[[i]][[ All ,1]], databg2[[i ]][[ All ,2]]\}]];
nlmbg2=Table[1,\{i,98\}];
nlmbg2 $=$ Table[NonlinearModelFit[databg2[[i]],A $* \operatorname{Exp}\left[(-1) *(1 / 2) *(x-\backslash[\mathrm{Mu}])^{\wedge} 2 *(1 / \backslash[\text { Sigma }])^{\wedge} 2\right],\{\{\backslash[$ Sigma
$], 0.05\},\{\backslash[\mathrm{Mu}], 19.4\},\{\mathrm{A}, 600\}\}, x],\{i, 98\}] ;$
 arbitrary units)" $\}$,ImageSize $->500$, PlotRange $->\{x, 16.7,17.95\}$, PlotLegends $->\{$ "With background intensity","Without background intensity"\}]Show[ListPlot[Thread[\{data2[[10]][[All,1]],data2[[[10]]][[All ,2]]-Mean[data2 [[10]][[1;;3,2]]]\}], PlotStyle->Red,PlotRange->All,ImageSize->500,AxesLabel->\{" $2 \backslash[$ Theta $]($ degree)", "Intensity (arbitrary units)" $\}], \operatorname{Plot}[n \operatorname{lm} 2[[10]][x],\{x, 19.25,19.55\}$, PlotRange $->$ All, Filling $->$ Bottom,ImageSize $->500$,AxesLabel $->\{" 2 \backslash[$ Theta](degree)","Intensity (arbitrary units)" $\}$ ]]

ListPlot[\{thickness1VsTime,thickness1CSVsTime\}, PlotRange->All,Joined-> True,AxesLabel->\{"Time ( minutes)","Thickness (nm)"\},PlotLegends->\{"Cubic crystallite","Spherical crystallite"\},Mesh->All, PlotMarkers->\{Automatic, 7 \},ImageSize->500]

ListPlot[\{thickness2VsTime,thickness2CSVsTime\},PlotRange->All,Joined->True,AxesLabel->\{"Time ( minutes)","IThickness (nm)"\},PlotLegends->\{"Cubic crystallite","Spherical crystallite"\},Mesh->All, PlotMarkers $->\{$ Automatic, 7$\}$,ImageSize $->500$ ]
ListPlot[\{thickness3VsTime,thickness3CSVsTime\}, PlotRange->All,Joined-> True,AxesLabel->\{"Time (
minutes)","Thickness (nm)"\},PlotLegends->\{"Cubic crystallite","Spherical crystallite"\},Mesh->All, PlotMarkers $->$ \{Automatic, 7 \},ImageSize - $>500$ ]

ListPlot [\{thickness4VsTime,thickness4CSVsTime\}, PlotRange->All,Joined->True,AxesLabel->\{"Time ( minutes)", "Thickness (nm)"\},PlotLegends->\{"Cubic crystallite", "Spherical crystallite"\},Mesh->All, PlotMarkers $->\{$ Automatic, 7$\}$,ImageSize $->500$ ]
ListPlot[\{thickness5VsTime,thickness5CSVsTime\},PlotRange $->$ All, Joined $->$ True,AxesLabel $->\{$ "Time ( minutes)","Thickness (nm)"\},PlotLegends->\{"Cubic crystallite","Spherical crystallite"\},Mesh->All, PlotMarkers->\{Automatic,7\},ImageSize->500]
(*Size estimation by IntegralBreadth method*)
IntegralBreadth1 =Table[1,\{i,98\}];
IntegralBreadth1 $=$ Table[Integrate[nlm1 $[[i]][x],\{x, 16.65,16.95\}] / n \operatorname{lm} 1[[i]][$ "ParameterTableEntries" ][[3,1]],\{ i ,98\}];
IntegralBreadth2 $=$ Table[1, $\{i, 98\}]$;
IntegralBreadth2 $=$ Table[Integrate[nlm2[[i ]][ $x],\{x, 19.25,19.55\}] / n \operatorname{lm} 2[[i]][$ "ParameterTableEntries"
][[3,1]],\{ i ,98\}];
IntegralBreadth3 =Table[1,\{i,98\}];
IntegralBreadth3 =Table[Integrate[nlm3[[i ]][x],\{x,27.35,27.75\}]/nlm3[[i ]][ "ParameterTableEntries" ][[3,1]],\{ i ,98\}];
IntegralBreadth4 $=$ Table[1,\{i,98\}];
IntegralBreadth4 $=$ Table[Integrate[nIm4[[i ]][x],\{x,32.25,32.65\}]/nlm4[[i ]][ "ParameterTableEntries" ][[3,1]],\{ i ,98\}];
IntegralBreadth5 $=$ Table $[1,\{\mathrm{i}, 98\}$ ];
IntegralBreadth5 $=$ Table[Integrate[nlm5[[i ]][ $x$ ], $\{x, 33.75,34.05\}] / n \operatorname{lm} 5[[i]][$ "ParameterTableEntries" ][[3,1]],\{ i ,98\}];

```
(*Find thickness *)
```

thicknessIB1 $=\left(1.1547 *\left(0.6941 / 10^{\wedge} 10\right)\right) /(($ IntegralBreadth $1 * P i / 180) * \operatorname{Cos}[$ re111[[2,All,1]]/2Degree $]) ;$
thicknessIB2 $=\left(1.0 *\left(0.6941 / 10^{\wedge} 10\right)\right) /(($ IntegralBreadth $2 * \mathbf{P i} / 180) * \operatorname{Cos}[$ re200 $[[2$, All, 1$]] / 2$ Degree $]) ;$
thicknessIB3 $=\left(1.0607 *\left(0.6941 / 10^{\wedge} 10\right)\right) /(($ IntegralBreadth $3 * \mathbf{P i} / 180) * \operatorname{Cos}[r e 220[[2$, All, 1$]] / 2$ Degree $]) ;$
thicknessIB4 $=\left(1.1359 *\left(0.6941 / 10^{\wedge} 10\right)\right) /(($ IntegralBreadth4*Pi/180 $) * \operatorname{Cos[re311[[2,All,1]]/2Degree]);~}$
thicknessIB5 $=\left(1.1547 *\left(0.6941 / 10^{\wedge} 10\right)\right) /(($ IntegralBreadth5 $* \mathbf{P i} / 180) * \operatorname{Cos}[r e 222[[2$, All,1]]/2Degree] $) ;$
thicknessIB1VsTime=Thread[\{time,thicknessIB1*10^9\}];
thicknessIB2VsTime $=$ Thread[\{time,thicknessIB2*10^9\}];
thicknessIB3VsTime $=$ Thread[\{time,thicknessIB3*10^9\}];
thicknessIB4VsTime=Thread[\{time,thicknessIB4*10^9\}];
128 thicknessIB5VsTime=Thread[\{time,thicknessIB5*10^9\}];

ListPlot[\{thickness5VsTime,thicknessIB5VsTime\}, PlotRange->All,Joined-> True,AxesLabel->\{"Time ( minutes)", "Thickness (nm)"\},PlotLegends->Placed[\{"FWHM","Integral breadth"\},\{0.85,0.2\}],Mesh->

All,PlotMarkers->\{Automatic,7\},ImageSize->500]
(*\{Comparing the estimated size growth rate of silver by FWHM and integral breadth method by shifting up the plot of size estimated by integral breadth method*)
thicknessIB1fVsTime=Thread[\{time,thicknessIB1*10^9\}];
thicknessIB2fVsTime=Thread[\{time,thicknessIB2*10^9\}];
thicknessIB3fVsTime=Thread[\{time,thicknessIB3*10^9\}];
thicknessIB4fVsTime=Thread[\{time,thicknessIB4*10^9\}];
thicknessIB5fVsTime=Thread[\{time,thicknessIB5*10^9\}];thickness1fVsTime=Thread[\{time,thickness1*10^9+ Mean[thicknessIB1*10^9-thickness $\left.\left.1 * 10^{\wedge} 9\right]\right\}$ ];
thickness2fVsTime=Thread[\{time,thickness $2 * 10^{\wedge} 9+$ Mean[thicknessIB2*10^9-thickness2*10^9]\}];
thickness3fVsTime=Thread[\{time,thickness $3 * 10^{\wedge} 9+$ Mean[thicknessIB3*10^9-thickness3*10^9]\}];
thickness4fVsTime=Thread[\{time,thickness $4 * 10^{\wedge} 9+$ Mean[thicknessIB4*10^9-thickness $\left.\left.4 * 10^{\wedge} 9\right]\right\}$ ];
thickness5fVsTime=Thread[\{time,thickness5*10^9+Mean[thicknessIB5*10^9-thickness5*10^9]\}];ListPlot[\{ thickness1fVsTime,thicknessIB1VsTime\}, PlotRange->All, Joined $->$ True, AxesLabel $->\{$ "Time (minutes) ","Thickness (nm)"\},PlotLegends->Placed[\{"FWHM","Integral breadth"\},\{0.85,0.35\}],Mesh->All, PlotMarkers $->\{$ Automatic, 7 \},ImageSize $->500$ ]
ListPlot[\{thickness2fVsTime, thicknessIB2VsTime\}, PlotRange->All, Joined-> True, AxesLabel->\{"Time ( minutes)","Thickness (nm)"\},PlotLegends->Placed[\{"FWHM","Integral breadth"\},\{0.85,0.35\}],Mesh-> All,PlotMarkers $->\{$ Automatic, 7$\}$,ImageSize->500]
ListPlot[\{thickness3flsTime, thicknessIB3VsTime\}, PlotRange->All,Joined->True,AxesLabel->\{"Time ( minutes)", "Thickness (nm)"\},PlotLegends->Placed[\{"FWHM","Integral breadth"\},\{0.85,0.35\}],Mesh-> All,PlotMarkers-> \{Automatic, 7 \},ImageSize->500]

ListPlot[\{thickness4fVsTime, thicknessIB4VsTime\}, PlotRange->All,Joined-> True,AxesLabel->\{"Time ( minutes)","Thickness (nm)"\},PlotLegends->Placed[\{"FWHM","Integral breadth"\},\{0.85,0.35\}],Mesh-> All,PlotMarkers-> \{Automatic, 7 \},ImageSize->500]
ListPlot[\{thickness5fVsTime, thicknessIB5VsTime\},PlotRange $->$ All,Joined $->$ True,AxesLabel $->\{$ "Time ( minutes)", "Thickness (nm)"\},PlotLegends->Placed[\{"FWHM","Integral breadth"\},\{0.85,0.35\}],Mesh-> All,PlotMarkers-> \{Automatic, 7 \}, ImageSize->500]
(*Size estimation error of FWHM and IntegralBreadth*)
nlm1bands95[x_]=Table[1,\{i,98\}];
nlm1bands95[x_]=Table[nlm1[[i]]]["MeanPredictionBands",ConfidenceLevel->.95],\{i,98\}];
nlm2bands95[x_]=Table[1,\{i,98\}];
nlm2bands95[x_]=Table[nlm2[[i]]["MeanPredictionBands",ConfidenceLevel->.95], \{i,98\}];
nlm3bands95[x_]=Table[1, $\{\mathrm{i}, 98\}]$;
nlm3bands95[x_]=Table[nlm3[[i]]["MeanPredictionBands",ConfidenceLevel->.95],\{i,98\}];
nlm4bands95[x_]=Table[1,\{i,98\}];
nlm4bands95[x_]=Table[nlm4[[i]]]["MeanPredictionBands",ConfidenceLevel->.95],\{i,98\}];
nlm5bands95[x_]=Table[1,\{i,98\}];
nlm5bands95[x_]=Table[nIm5[[i]]["MeanPredictionBands",ConfidenceLevel->.95],\{i,98\}];
IntergralBreadth1delta $=$ Table $[1,\{i, 98\}] ;$
IntergralBreadth1delta $=$ Table[(Abs[Integrate[nlm1bands95[x][[i, 1]],\{ $x, 16.65,16.95\}] / n \operatorname{lm} 1[[i]][$ "
ParameterConfidenceIntervals " , ConfidenceLevel - >.95][[3,2]]- IntegralBreadth1 [[ i $]$ ]]+ Abs[Integrate[
nlm1bands95[x][[i ,2]],\{ x ,16.65,16.95\}]/nlm1[[i ]][ " ParameterConfidenceIntervals ", ConfidenceLevel
->.95][[3,1]]- IntegralBreadth1 [[ i ]]]) /2,\{i ,98\}];
IntergralBreadth2delta =Table[1,\{i,98\}];
IntergralBreadth2delta $=$ Table[(Abs[Integrate[nlm2bands95[x][[i, 1]],\{ x , 19.25, 19.55\}]/ nlm2[[i ]][ "
ParameterConfidenceIntervals " , ConfidenceLevel ->.95][[3,2]]- IntegralBreadth2 [[ i $]$ ]] + Abs[Integrate[
nlm2bands95[x][[i ,2]],\{ x , 19.25,19.55\}]/nIm2[[i ]][ " ParameterConfidenceIntervals " , ConfidenceLevel
->.95][[3,1]]- IntegralBreadth2 [[ i ]]]) /2,\{i ,98\}];
IntergralBreadth3delta $=$ Table[1,\{i,98\}];
IntergralBreadth3delta $=$ Table[(Abs[Integrate[nlm3bands95[x][[i ,1]],\{ $x, 27.35,27.75\}] / n \operatorname{lm} 3[[\mathrm{i}]][$ "
ParameterConfidenceIntervals " , ConfidenceLevel ->.95][[3,2]]- IntegralBreadth3 [[ i ]]]+Abs[Integrate[
nlm3bands95[x][[i ,2]],\{ x,27.35,27.75\}]/nlm3[[i ]][ " ParameterConfidenceIntervals", ConfidenceLevel
->.95][[3,1]]- IntegralBreadth3 [[i ]]]) /2, \{i ,98\}];
IntergralBreadth4delta $=$ Table[1,\{i,98\}];
IntergralBreadth4delta $=$ Table[(Abs[Integrate[nlm4bands95[x][[i,1]],\{x,32.25,32.65\}]/nlm4[[i ]][ "
ParameterConfidenceIntervals " , ConfidenceLevel - > 95][[3,2]]- IntegralBreadth4 [[ i $]$ ]]+Abs[Integrate[
nlm4bands95[x][[i ,2]],\{ x , 32.25,32.65\}]/nlm4[[i ]][ " ParameterConfidenceIntervals ", ConfidenceLevel
->.95][[3,1]]- IntegralBreadth4 [[ i ]]]) /2, \{ i ,98\}];
IntergralBreadth5delta $=$ Table[1,\{i,98\}];
IntergralBreadth5delta $=$ Table[(Abs[Integrate[nIm5bands95[x][[i,1]],\{ x, 33.75,34.05\}]/ nlm5[[i ]][ "
ParameterConfidenceIntervals " , ConfidenceLevel ->.95][[3,2]]- IntegralBreadth5 [[ i ]]] + Abs[Integrate[ nlm5bands95[x][[i ,2]],\{ x , 33.75,34.05\}]/ nlm5[[i ]][ " ParameterConfidenceIntervals ", ConfidenceLevel ->.95][[3,1]]- IntegralBreadth5 [[ i ]]]) /2, \{ i ,98\}]; FWHM1delta=Table[1, \{i,98\}];

FWHM2delta=Table[1,\{i,98\}];
FWHM3delta=Table[1,\{i,98\}];
FWHM4delta=Table[1,\{i,98\}];
FWHM5delta=Table[1,\{i,98\}];
FWHM1delta = Table[(Abs[nlm1[[i]]["ParameterConfidenceIntervals",ConfidenceLevel->.95][[1,1]]*2*Sqrt[2* $\boldsymbol{L o g}[2]]-$ FWHM1[[i]]] $+\mathbf{A b s [ n I m 1 [ [ i ] ] [ " P a r a m e t e r C o n f i d e n c e I n t e r v a l s " , ~ C o n f i d e n c e L e v e l ~ - ~ > . 9 5 ] [ [ 1 , 2 ] ] ] * 2 * S q r t ~}$
[2*Log[2]]-FWHM1[[i]]])/2,\{i,98\}];
${ }^{187}$ FWHM2delta= Table[(Abs[nlm2[[i]]]["ParameterConfidenceIntervals",ConfidenceLevel->.95][[1,1]]*2*Sqrt[2* $\log [2]]-$ FWHM2[[i]]]]+Abs[nlm2[[i]]["ParameterConfidencelntervals",ConfidenceLevel ->.95][[1,2]]*2*Sqrt [2*Log[2]]-FWHM2[[i]]])/2,\{i,98\}];
${ }_{188}$ FWHM3delta= Table[(Abs[nlm3[[i]]["ParameterConfidencelntervals",ConfidenceLevel->.95][[1,1]]*2*Sqrt[2* $\log [2]]-$ FWHM3[[i]]] $+\mathbf{A b s [ n l m 3 [ [ i ] ] [ " P a r a m e t e r C o n f i d e n c e I n t e r v a l s " , C o n f i d e n c e L e v e l ~ - ~}-95][[1,2]] * 2 *$ Sqrt [2*Log[2]]-FWHM3[[i]]])/2,\{i,98\}];
FWHM4delta= Table[(Abs[nlm4[[i]]]["ParameterConfidencelntervals",ConfidenceLevel->.95][[1,1]]*2*Sqrt[2* $\log [2]]-$ FWHM4[[i]]] $+\mathbf{A b s [ n l m 4 [ [ i ] ] [ " P a r a m e t e r C o n f i d e n c e I n t e r v a l s " , C o n f i d e n c e L e v e l ~ - > . 9 5 ] [ [ 1 , 2 ] ] * 2 * S q r t ~}$ [2*Log[2]]-FWHM4[[i]]])/2,\{i,98\}];
$204\left(\left(\left(0.6941 / 10^{\wedge} 10\right) * 0.2138\right) /((F W H M 3 * P i / 180) *((\operatorname{Cos}[r e 220[[2, \text { All,1] }] / 2 D e g r e e]))))^{\wedge} 2+\right.$
$205(((0.8551) *(0.000000000000006941)) /((F W H M 3 * P i / 180) *(\operatorname{Cos}[r e 220[[2$, All,1] $] / 2 D e g r e e]))) \wedge 2+$
${ }_{206}\left(\left((0.8551) *\left(0.6941 / 10^{\wedge} 10\right) *(\mathrm{FWHM} 3 d e l t a * P i / 180)\right) /\left((\mathrm{FWHM} 3 * \mathrm{Pi} / 180)^{\wedge} 2 *(\operatorname{Cos}[\mathrm{re} 220[[2, \mathbf{A l l}, 1]] / 2\right.\right.$
Degree]) )) ${ }^{\wedge} 2+$
( ( (0.8551) *(0.6941/10^10) * (Sin[re220 [[2, All ,1]]/2 Degree]) $*(0.034842237140) ~) /(($ FWHM3*Pi/180) $)$
$\left.\left(\operatorname{Cos}[r e 220[[2, A l l, 1]] / 2 \text { Degree] })^{\wedge} 2\right)\right)^{\wedge} 2$ ];deltathicknessFWHM4 $=$ Sqrt [
$208\left(\left(\left(0.6941 / 10^{\wedge} 10\right) * 0.0426\right) /((F W H M 4 * P i / 180) *((\operatorname{Cos}[r e 311[[2, A l l, 1]] / 2 D e g r e e])))\right)^{\wedge} 2+$
$209\left(((0.8551) *(0.000000000000006941)) /((F W H M 4 * P i / 180) *(\operatorname{Cos}[\text { re311[[2,All,1]]/2Degree] }) \quad))^{\wedge} 2+\right.$ ${ }_{210}\left(\left((0.8551) *\left(0.6941 / 10^{\wedge} 10\right) *(\right.\right.$ FWHM4delta $\left.* P i / 180)\right) /(($ FWHM4*Pi/180) $) 2 *(\operatorname{Cos}[r e 311[[2, A l l, 1]] / 2$ Degree]) ) $)^{\wedge} 2+$
$\left(\left((0.8551) *\left(0.6941 / 10^{\wedge} 10\right) *(\operatorname{Sin}[r e 311[[2\right.\right.$, All , 1] $] / 2$ Degree $]) *(0.034842237140)) /(($ FWHM4*Pi/180 $) *$ $\left.\left.(\boldsymbol{C o s}[r e 311[[2, \text { All }, 1]] / 2 \text { Degree }])^{\wedge} 2\right)\right)^{\wedge} 2$ ];deltathicknessFWHM5 = Sqrt [
$\left(\left(\left(0.6941 / 10^{\wedge} 10\right) * 0.0849\right) /((\text { FWHM5 } * \text { Pi/180 }) *((\operatorname{Cos}[\text { re222[[2,All,1]]/2Degree }])))\right)^{\wedge} 2+$
$(\quad((0.8551) *(0.000000000000006941)) /((F W H M 5 * P i / 180) *(\operatorname{Cos}[r e 222[[2, A l l, 1]] / 2 D e g r e e]) \quad))^{\wedge} 2+$ $\left(\left((0.8551) *\left(0.6941 / 10^{\wedge} 10\right) *(\text { FWHM5delta*Pi/180) ) / ( (FWHM5*Pi/180) })^{2} 2 *(\operatorname{Cos}[r e 222[[2, A l l, 1]] / 2\right.\right.$ Degree]) ) ) ^2 +
$\left(\left((0.8551) *\left(0.6941 / 10^{\wedge} 10\right) *(\operatorname{Sin}[r e 222[[2\right.\right.$, All ,1] $] / 2$ Degree $]) *(0.034842237140)) /(($ FWHM5 $5 \mathrm{Pi} / 180) *$ $\left(\operatorname{Cos}\left[r e 222[[2, \text { All, 1] ]/2Degree }])^{\wedge} 2\right)\right)^{\wedge} 2$ ];
( $\quad((0.8551) *(0.000000000000006941)) /(($ IntegralBreadth $3 * P i / 180) *(\operatorname{Cos}[r e 220[[2$, All ,1]]/2 Degree] $)$ ) )^2+
${ }_{231}\left(\left((0.8551) *\left(0.6941 / 10^{\wedge} 10\right) *(\right.\right.$ IntergralBreadth3delta $\left.* P i / 180)\right) / \quad(($ IntegralBreadth3 $* P i / 180) \wedge 2 * \quad($ $\operatorname{Cos}\left[r e 220[[2, \text { All ,1]]/2 Degree] }) \text { ) })^{\wedge} 2+\right.$
$232\left(\left((0.8551) *\left(0.6941 / 10^{\wedge} 10\right) *(\operatorname{Sin}[r e 220[[2\right.\right.$, All ,1] $] / 2$ Degree $]) *(0.034842237140)) /(($ IntegralBreadth3 $*$ Pi/180)* $\left(\mathbf{C o s}\left[r e 220\left[[2, \text { All ,1]]/2 Degree] })^{\wedge} 2\right)\right)^{\wedge} 2\right.$ ]; deltathicknessIntegralBreadth4 = Sqrt [
$(\quad((0.8551) *(0.000000000000006941)) /(($ IntergralBreadth5delta $* P i / 180) *(\mathbf{C o s}[r e 222[[2$, All ,1]]/2 Degree] ) ) ^^2+

ListPlot $\left[\left\{\right.\right.$ thickness $3 * 10^{\wedge} 9$, thickness $3 * 10^{\wedge} 9+$ deltathicknessFWHM3 $* 10^{\wedge} 9$, thickness $3 * 10^{\wedge} 9-$ deltathicknessFWHM3*10^9\},Joined->True,PlotRange->All,PlotLegends->Placed[\{"size estiamtion"," upper limit of size estimation", "lower limit of size estimation" $\},\{0.7,0.2\}]$, AxesLabel-> "Time (minutes) ","Thickness (nm)"\},ImageSize->500]
ListPlot $\left[\left\{\right.\right.$ thickness $4 * 10^{\wedge} 9$, thickness $4 * 10^{\wedge} 9+$ deltathicknessFWHM4*10^9,thickness $4 * 10^{\wedge} 9-$
deltathicknessFWHM4*10^9\},Joined->True,PlotRange->All,PlotLegends->Placed[\{"size estiamtion"," upper limit of size estimation", "lower limit of size estimation" $\},\{0.7,0.2\}]$, AxesLabel-> "Time (minutes) ","Thickness (nm)"\},ImageSize->500]

ListPlot[\{thicknessIB5*10^9,thicknessIB5*10^9+deltathicknessIntegralBreadth5*10^9,thicknessIB5*10^9-deltathicknessIntegralBreadth5*10^9\},Joined->True,PlotRange->All,PlotLegends->Placed[\{"size estiamtion","upper limit of size estimation","lower limit of size estimation" $\},\{0.7,0.2\}]$, AxesLabel-> \{ "Time (minutes)", "thicknessIB (nm)"\},ImageSize->500]

## Listing D.5: Chapter 5

(*to show the jumping of diffraction angle due to insufficient diffration angle resolution *)

ListPlot[sample [[2,5;;8]], PlotRange->\{\{19.25,19.55\},\{20,120\}\},Joined->True,Mesh->All,ImageSize ->500,PlotMarkers->Automatic,PlotLegends->\{"388\[Degree]C","390\[Degree]C","392\[Degree]C"," 394\[Degree]C"\},AxesLabel-> \{"2\[Theta](degree)","Intensity (arbitrary units)"\}]
(*Comparing theoretical and experiment values of diffraction angle of relfection 200*)

```
theoreticalExpo2=Flatten[Import["C:\\Users\\roy\\\Desktop\\thesis\\theoretical diffraction angle - 200
    expo.xlsx" ]];
theoreticalExpoc2=Flatten[Import["C:\\\Users\\roy\\\Desktop\\thesis\\theoretical diffraction angle - 200
    expoc.xlsx" ]];
ListPlot[{Thread[{Temp,theoreticalExpoc2}],Thread[{Temp,re200[[2,AlI,1]]}]},PlotLegends-> {"Theoretical",
    "Experimental"},AxesLabel-> {"Temperature(\[Degree]C)","2\[Theta](degree)"},Joined-> True,
    ImageSize->500,PlotMarkers->{"",{\[FilledSmallCircle],10}}]
(*Comparing thickness with and without energy dispersion broadening for reflection 111*)
FWHMEnergyDispersion=Flatten[Import["C:\\Users\\roy\\Desktop\\thesis\\\broadening due to energy
    dispersion result(111).xlsx"]];
FWHMWithoutEnergyDispersion=Sqrt[FWHM1^2-FWHMEnergyDispersion^2];
thicknessWithoutEnergyDispersion=(0.8551*(0.6941/10^10))/((FWHMWithoutEnergyDispersion*Pi/180)*Cos[
    re111[[2,All,1]]/2Degree]);
thicknessWithoutEnergyDispersionVsTime=Thread[{time,thicknessWithoutEnergyDispersion*10^9}];
ListPlot[{thickness1VsTime,thicknessWithoutEnergyDispersionVsTime},PlotRange->All,Joined-> True,
    Mesh->All,ImageSize->500,PlotMarkers->{Automatic,7},AxesLabel->{"Time (minutes)","
    Thickness (nm)"},PlotLegends->Placed[{"thickness(with energy dispersion broadening)","thickness(
    without energy dispersion broadening)"},{0.66,0.3}]]
```

Listing D.6: Chapter 6

```
(*lsothermal kinetics *)
yfwhm01=Table[1,{i,25}];
yfwhm01=Table[thickness1[[i]]^3/thickness1 [[27]]^3,{ i ,2,26}];
yfwhm1=Log[-Log[1-yfwhm01]];
xfwhm1=Log[time[[2;;26]]*60];
yvsxfwhm1=Thread[{xfwhm1,yfwhm1}];
yfwhm02=Table[1,{i,25}];
yfwhm02=Table[thickness2[[i]]^3/thickness2 [[27]]^3,{ i ,2,26}];
yfwhm2=Log[-Log[1-yfwhm02]];
xfwhm2=Log[time[[2;;26]]*60];
yvsxfwhm2=Thread[{xfwhm2,yfwhm2}];
yfwhm03=Table[1,{i,25}];
```

```
yfwhm03=Table[thickness3[[i]]^3/thickness3 [[27]]^3,{ i ,2,26}];
yfwhm3=Log[-Log[1-yfwhm03]];
xfwhm3=Log[time[[2;;26]]*60];
yvsxfwhm3=Thread[{xfwhm3,yfwhm3}];
yfwhm04=Table[1,{i,25}];
yfwhm04=Table[thickness1[[i]]^3/thickness1[[27]]^3,{ i ,2,26}];
yfwhm4=Log[-Log[1-yfwhm01]];
xfwhm4=Log[time[[2;;26]]*60];
yvsxfwhm4=Thread[{xfwhm1,yfwhm1}];
yfwhm05=Table[1,{i,25}];
yfwhm05=Table[thickness5[[i]]^3/thickness5 [[27]]^3,{ i ,2,26}];
yfwhm5=Log[-Log[1-yfwhm05]];
xfwhm5=Log[time[[2;;26]]*60];
yvsxfwhm5=Thread[{xfwhm5,yfwhm5}];
yIB01=Table[1,{i,25}];
yIB01=Table[thicknessIB1[[i]]^3/thicknessIB1 [[27]]^3,{ i ,2,26}];
ylB1=Log[-Log[1-yIB01]];
xIB1=Log[time [[2;;26]]*60];
yvsxIB1=Thread[{xIB1,yIB1}];
ylB02=Table[1,{i,25}];
yIB02=Table[thicknessIB2[[i]]^3/thicknessIB2 [[27]]^3,{ i ,2,26}];
yIB2=Log[-Log[1-yIB02]];
xIB2=Log[time [[2;;26]]*60];
yvsxIB2=Thread[{xIB2,yIB2}];
yIB03=Table[1,{i,25}];
ylB03=Table[thicknessIB3[[i]]^3/ thicknessIB3 [[27]]^3,{ i ,2,26}];
ylB3=Log[-Log[1-ylB03]];
xIB3=Log[time [[2;;26]]*60];
yvsxIB3=Thread[{xIB3,yIB3}];
yIB04=Table[1,{i,25}];
```

```
ylB04=Table[thicknessIB1[[i]]^3/thicknessIB1 [[27]]^3,{ i ,2,26}];
yIB4=Log[-Log[1-yIB01]];
xIB4=Log[time [[2;;26]]*60];
yvsxIB4=Thread[{x|B1,ylB1}];
yIB05=Table[1,{i,25}];
yIB05=Table[thicknessIB5[[i]]^3/thicknessIB5 [[27]]^3,{ i ,2,26}];
yIB5=Log[-Log[1-yIB05]];
xIB5=Log[time [[2;;26]]*60];
yvsxIB5=Thread[{xlB5,ylB5}];
ListPlot[yvsxfwhm1,PlotRange->All,AxesLabel-> {"ln(t)"," |n[-ln(1-x(t))]"},ImageSize->500]
ListPlot[yvsxfwhm2,PlotRange->All,AxesLabel-> {" ln (t)"," ln[-ln(1-x(t))]"},ImageSize->500]
ListPlot[yvsxfwhm3,PlotRange->All,AxesLabel->{"In(t)","In[-\operatorname{ln}(1-x(t))]"},ImageSize->500]
ListPlot[yvsxfwhm4,PlotRange->AII,AxesLabel->{"ln(t)","ln[-In(1-x(t))]"},ImageSize->500]
ListPlot[yvsxfwhm5,PlotRange->AII,AxesLabel->{"ln(t)","ln[-ln(1-x(t))]"},ImageSize->500]
ListPlot[yvsx|B1,PlotRange->All,AxesLabel-> {"ln(t)","ln[-ln(1-x(t))]"},ImageSize->500]
ListPlot[yvsx|B2,PlotRange->AII,AxesLabel->{"ln(t)","ln[-\operatorname{ln}(1-x(t))]"},ImageSize->500]
ListPlot[yvsxIB3,PlotRange->AII,AxesLabel->{"ln(t)","ln[-ln(1-x(t))]"},ImageSize->500]
ListPlot[yvsxIB4,PlotRange->AII,AxesLabel->{"ln(t)","ln[-ln(1-x(t))]"},ImageSize->500]
ListPlot[yvsxIB5,PlotRange->AII,AxesLabel->{"ln(t)","ln[-In(1-x(t))]"},ImageSize->500]
Show[Graphics3D[Ellipsoid [{0,0,0},{92.37',76.71',86.57'}]], ViewPoint-> {1.3',-2.4',2.'}]
Show[Graphics3D[Ellipsoid [{0,0,0},{72.82',72.34',72.45'}],{ ViewPoint-> {1.3',-2.4',2.'}}],ViewPoint
    ->{1.3',-2.4',2.'}]
```


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[^0]:    ${ }^{1}$ See Appendix A for the calculated diffraction angles by linear regression and exponential regression

[^1]:    ${ }^{2}$ See Appendix A for the calculated diffraction angles by linear regression and exponential regression

[^2]:    ${ }^{3}$ See Appendix D for the code to avoid recording the side peak.

[^3]:    ${ }^{4}$ See Appendix D for the Mathematica code.

[^4]:    ${ }^{5}$ See Appendix B for the calculated scaling factors

[^5]:    ${ }^{1}$ See Appendix D for the Mathematica code.

[^6]:    ${ }^{1}$ See Appendix C for errors at other temperatures

